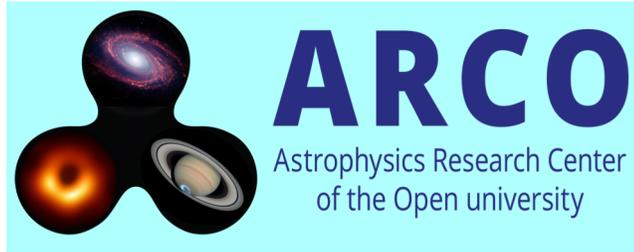


Short Gamma-Ray Bursts in the Era of Multi-Messenger Astrophysics



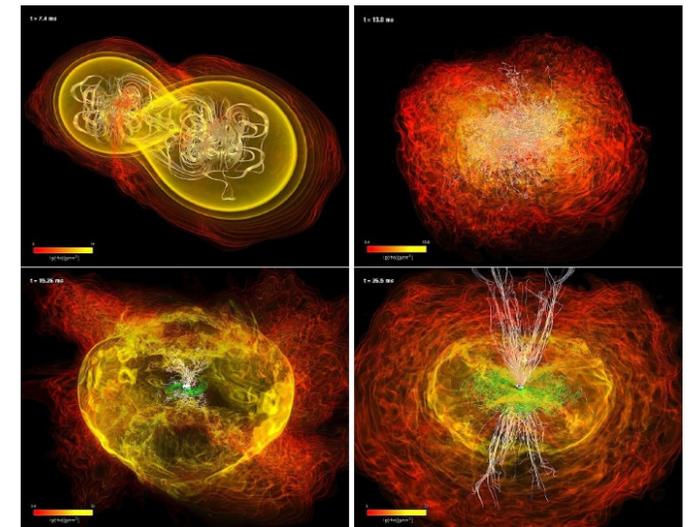
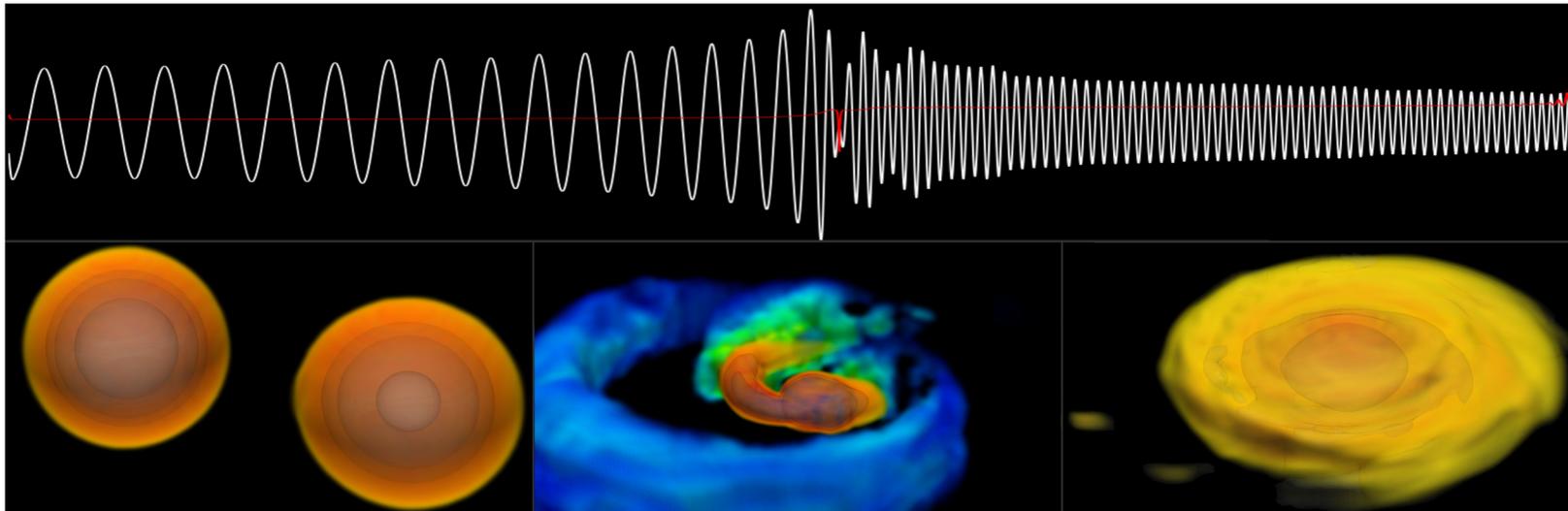
**Jonathan
Granot**



THE GEORGE
WASHINGTON
UNIVERSITY
WASHINGTON DC

Open University of Israel

& George Washington University



High Energy and the Cosmos; AGASS-HEP-Cosmology Workshop

Ariel University, 20 February 2025

GRB Theoretical Framework:

■ Progenitors:

❖ **Long:** massive stars

❖ **Short:** binary mergers (NS-NS, BH-NS?)

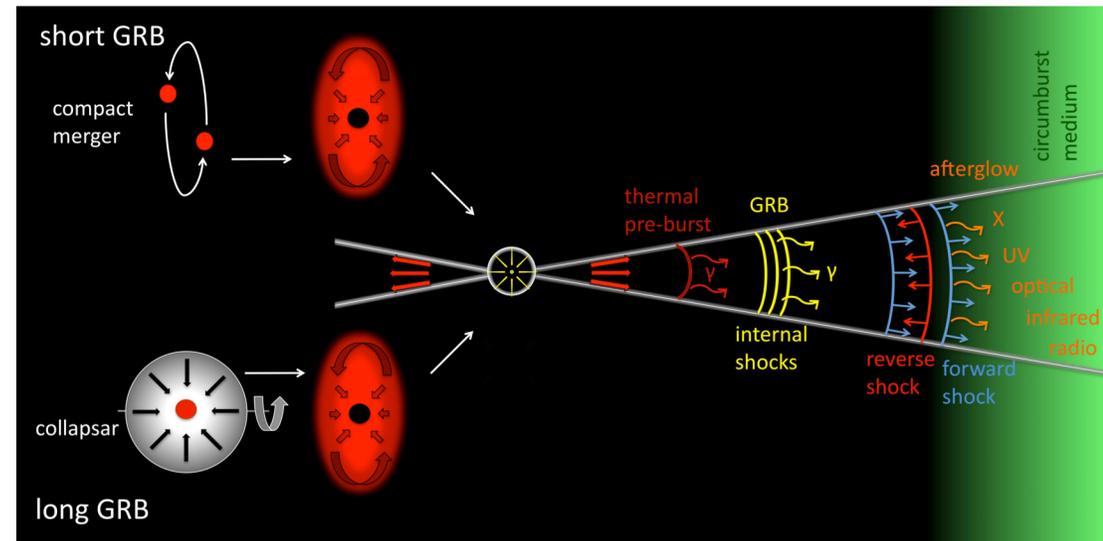
■ Acceleration: fireball or magnetic?

■ Prompt γ -rays: dissipation – internal shocks or magnetic reconnection?

Emission mechanism?

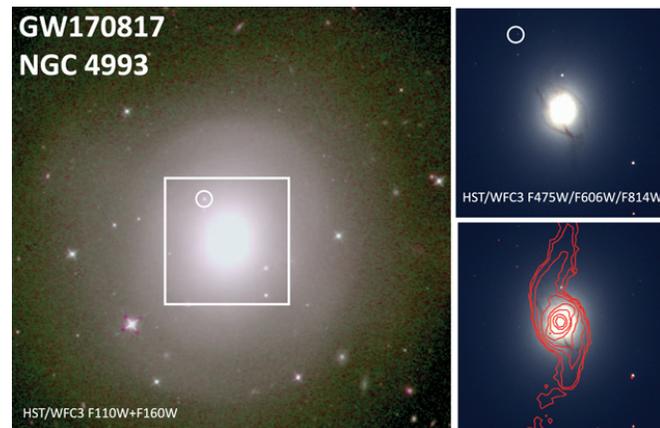
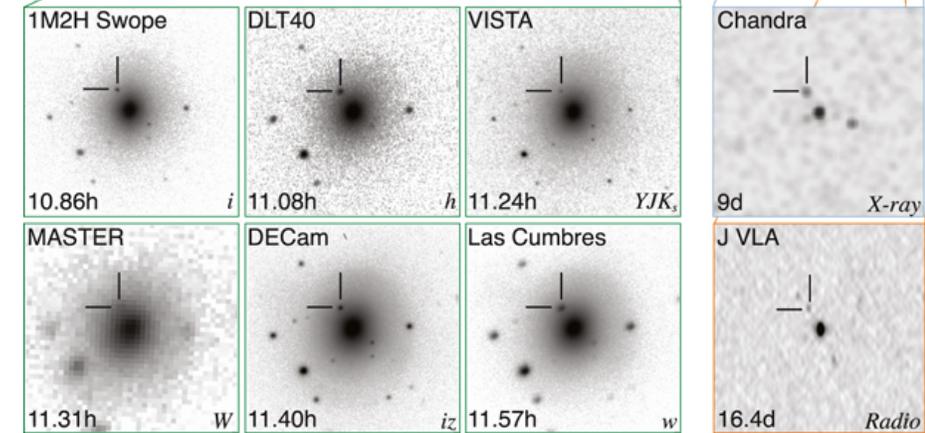
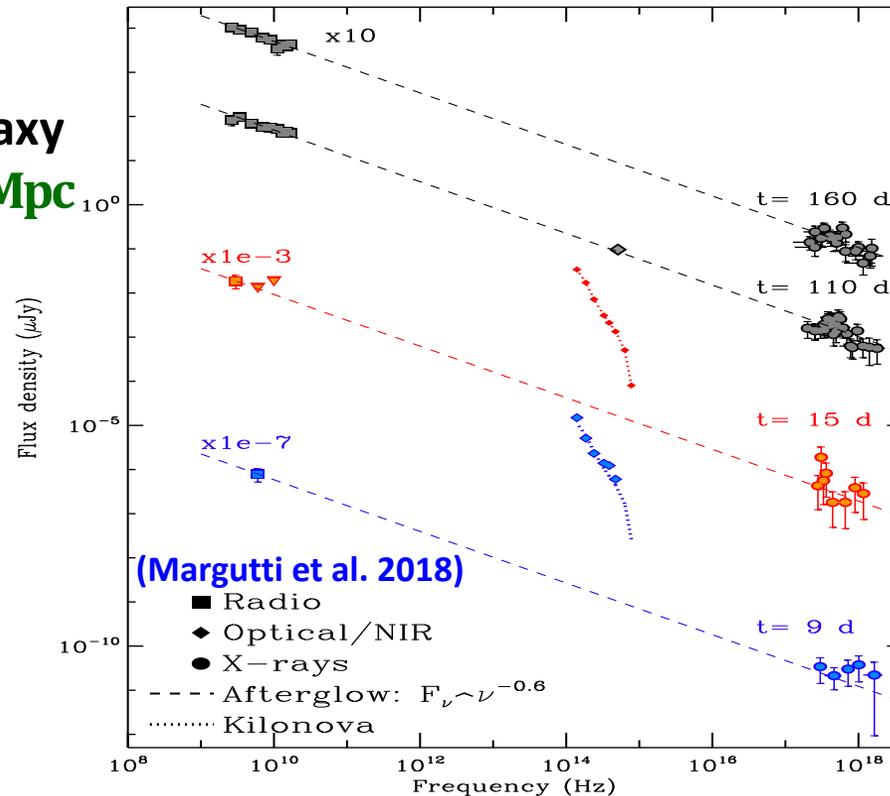
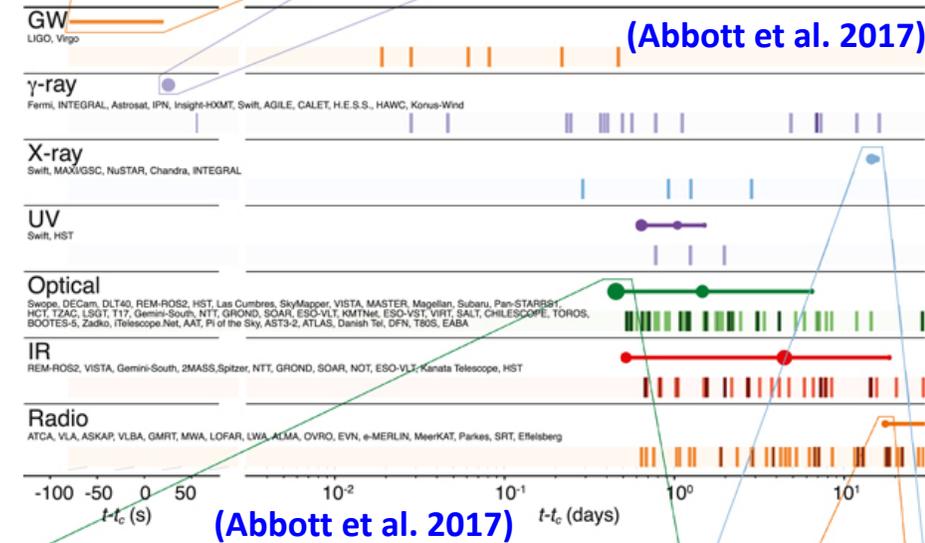
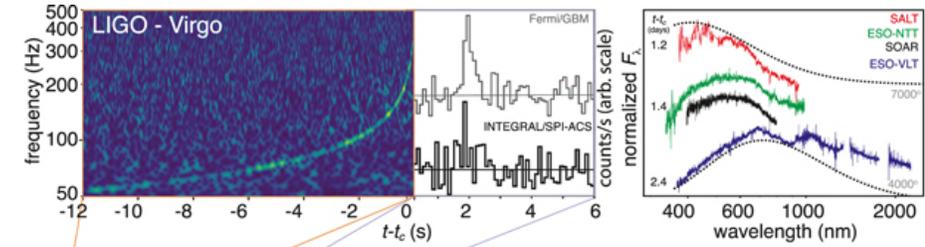
■ Deceleration: the outflow decelerates (by a reverse shock for $\sigma \lesssim 1$) as it sweeps-up the external medium

■ Afterglow: from the long lived forward shock going into the external medium; as the shock decelerates the typical frequency decreases: X-ray \rightarrow optical \rightarrow radio



GW170817 / GRB170817A: NS-NS merger

- First NS-NS merger detected in gravitational waves (GW)
- First electromagnetic counterpart to a GW event
 - The short GRB 170817A (very under-luminous, 1.74 s γ -GW delay)
 - Optical (IR to UV) kilonova emission over a few weeks
 - X-ray (> 9 d; still barely detected) to radio (>16 d) afterglow
- First direct sGRB - NS-NS merger association (Eichler+ 1989)
- First clear-cut kilonova
- $D_{GW} = 43_{-6.9}^{+2.9}$ Mpc; host galaxy is elliptical: $D = 41.0 \pm 3.1$ Mpc ($z = 0.009783$) 2 kpc from host center in projection



GW170817 / GRB170817A: Kilonova

Observations require two components:

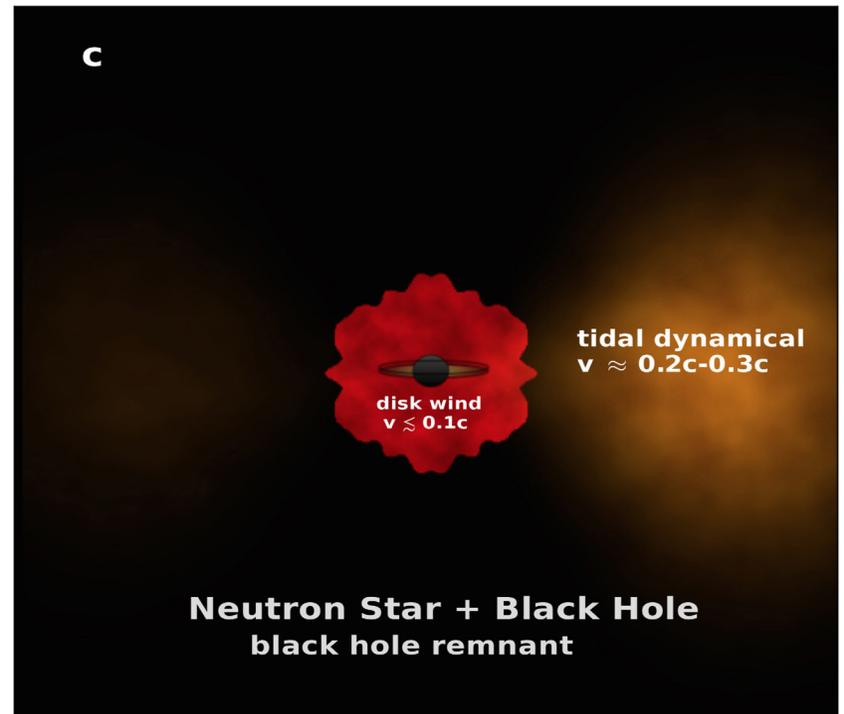
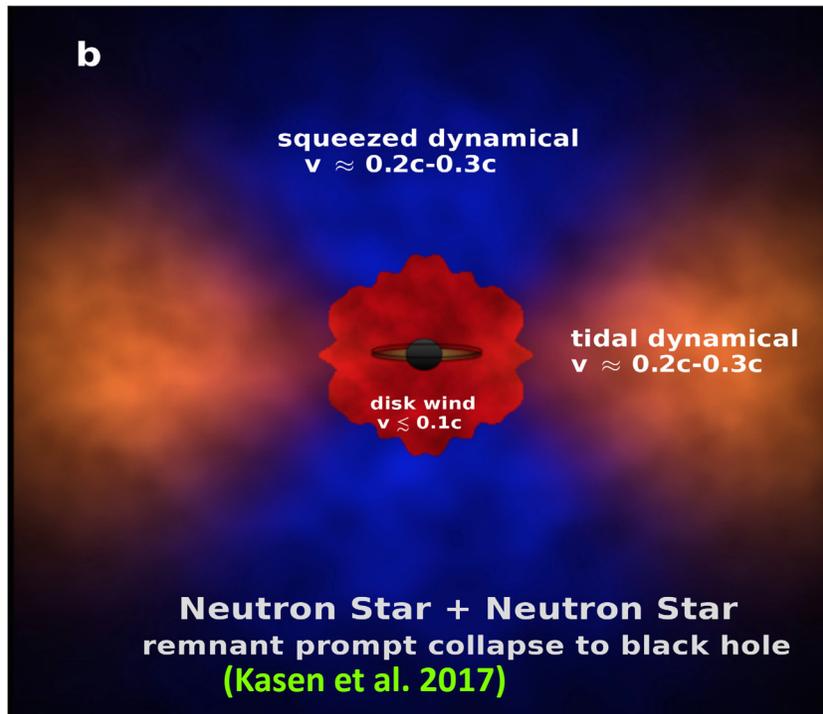
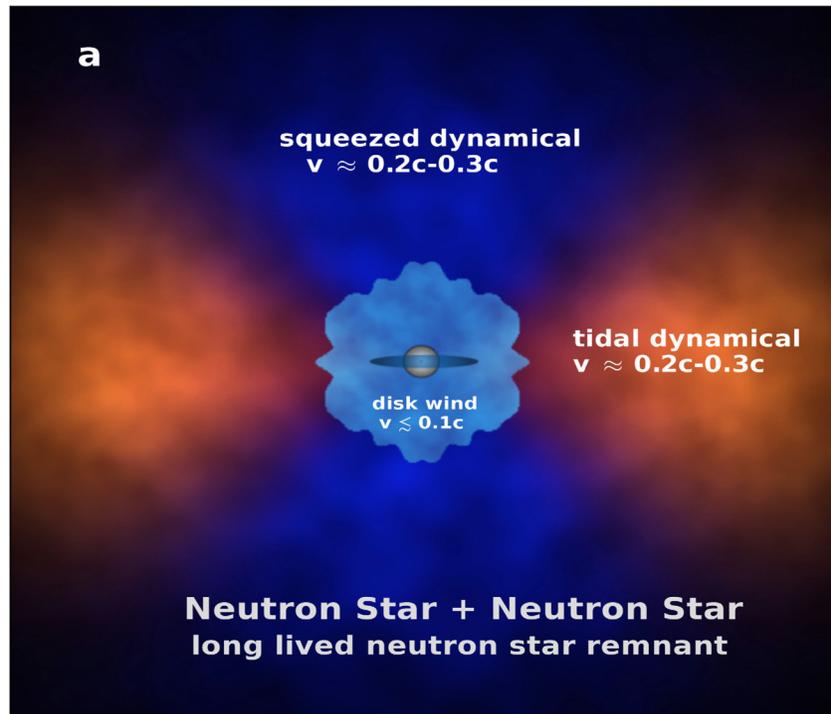
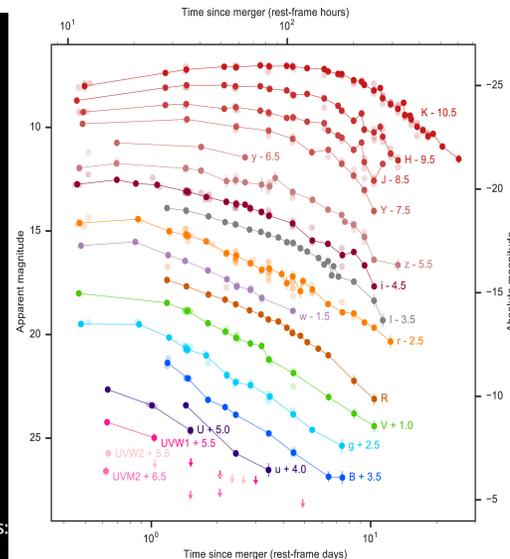
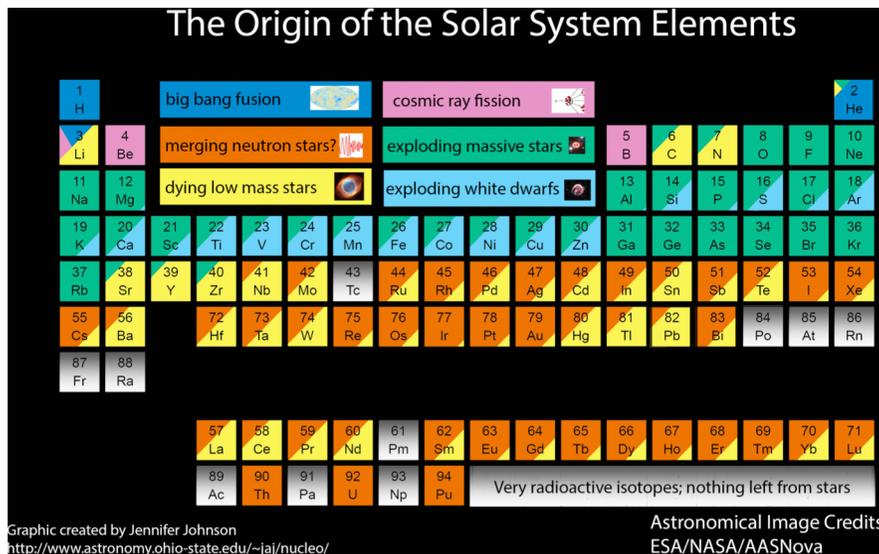
❖ First **blue/fast**, **lanthanide-poor**

$$M_{ej} \approx (1\% - 2\%)M_{\odot}, v_{ej} \approx (0.2 - 0.3)c$$

❖ Second **red/slow**, **lanthanide-rich**

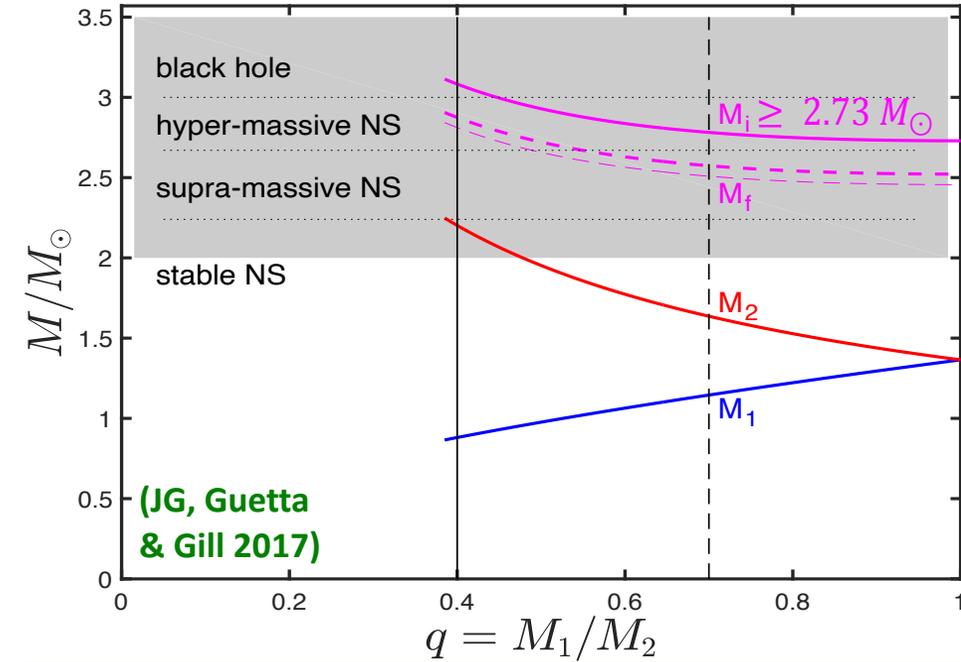
$$M_{ej} \approx (3\% - 5\%)M_{\odot}, v_{ej} \approx (0.05 - 0.2)c$$

Synthesized large amounts of heavy elements (may dominate the cosmic r-process nucleosynthesis, heavy metals e.g. gold, platinum)

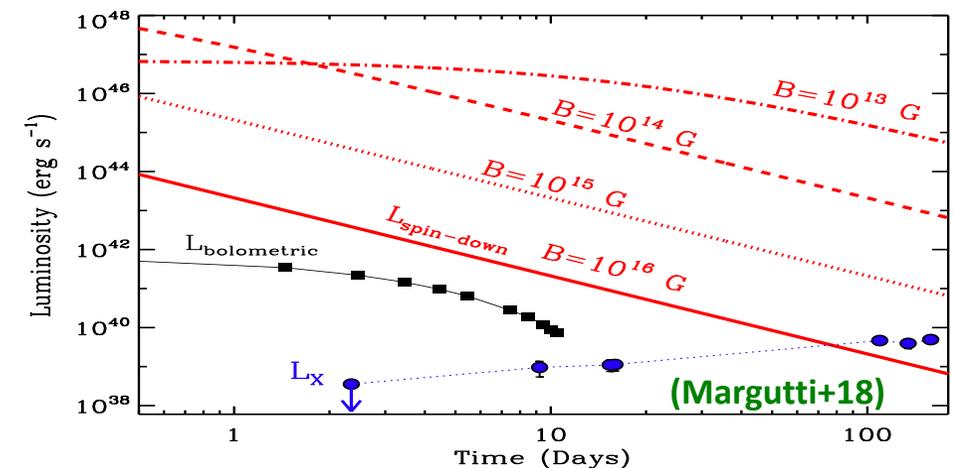


GW170817 / GRB170817A: Remnant Type

- $M_{1,2}$ = pre-merger NS $M_{\text{gravitational}}$
- post-merger total mass: $M_i = M_1 + M_2$
- Final mass $M_f \approx 0.93M_i$ due to:
 - ❖ GW & neutrino energy losses
 - ❖ Mass ejection during the merger
- A stable NS or SMNS $\Rightarrow P_0 \approx 1 \text{ ms} \Rightarrow E_{\text{rot}} \gtrsim 10^{52.5} \text{ erg}$,
 $\tau_{\text{sd}} \approx 20B_{13}^{-2} \text{ days} \Rightarrow$ would contradict afterglow observations (also what produces the GRB/afterglow?)
- The argument can be reversed to constrain NS EoS & $M_{\text{max}} \lesssim 2.17M_{\odot}$ (Margalit & Metzger 2017; Rezzolla et al. 2018)

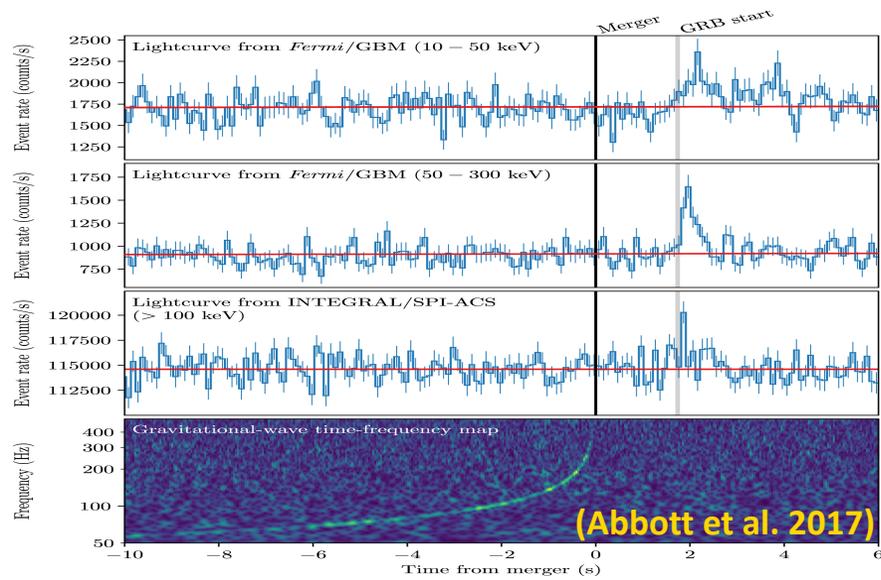


Chirp mass: $\mathcal{M} = \left(\frac{M_1^3 M_2^3}{M_1 + M_2} \right)^{1/5} = 1.188^{+0.004}_{-0.002} M_{\odot}$ (Abbott+ 2017)



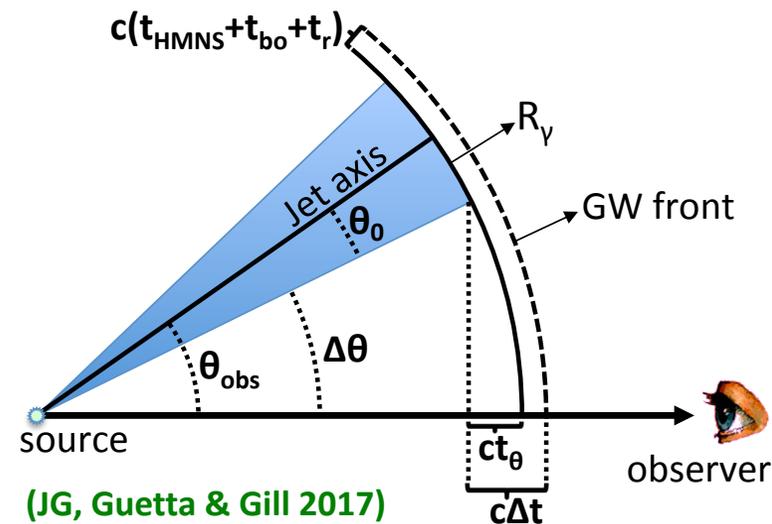
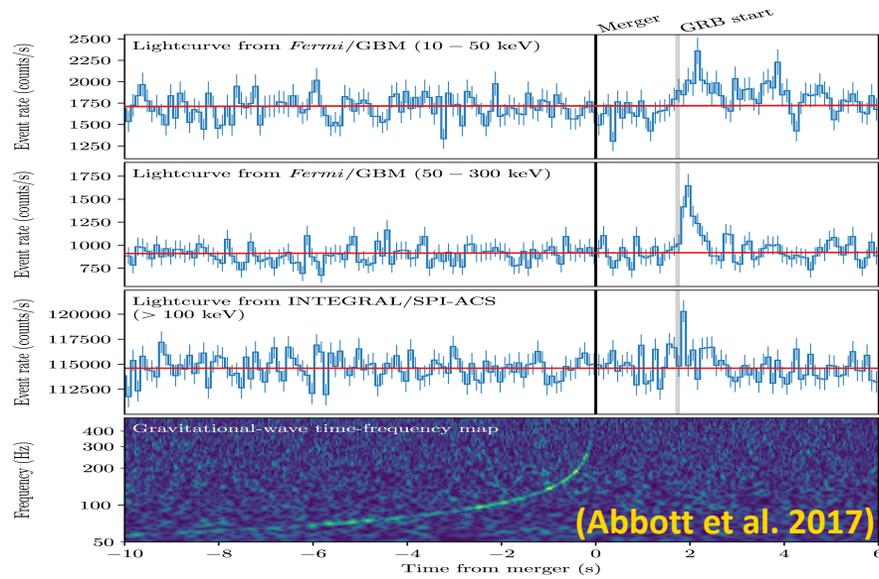
GW170817 / GRB170817A: The Time Delay

- The $\Delta t \approx 1.74$ s delay between the GW chirp signal & the sGRB onset $\Rightarrow \left| \frac{v_{\text{GW}}}{c} - 1 \right| \lesssim 4 \cdot 10^{-16}$
- A HMNS may explain $\Delta t \approx 1.74$ s by $t_{\text{HMNS}} \lesssim 0.5$ s & $t_{\text{bo}} \sim 1$ s
(Moharana & Piran 2017 find $t_{\text{bo}} \sim 0.5$ s for SGRBs, from a plateau in their duration distribution, $dN_{\text{GRB}}/dT_{\text{GRB}}$)
- Direct BH formation \Rightarrow a shorter jet breakout time $t_{\text{bo}} \Rightarrow$ the jet is less likely to be choked
- If the prompt γ -rays are beamed away from us (large $\Gamma\Delta\theta$), the implied on-axis $L_{\gamma,\text{iso}}$ & E_{peak} are very high – inconsistent with their observed correlation (JG+ 2017) & implying large compactness (Matsumoto+ 2019) \Rightarrow they must arise from $\Gamma\Delta\theta < 1 \Rightarrow$ a jet with angular structure



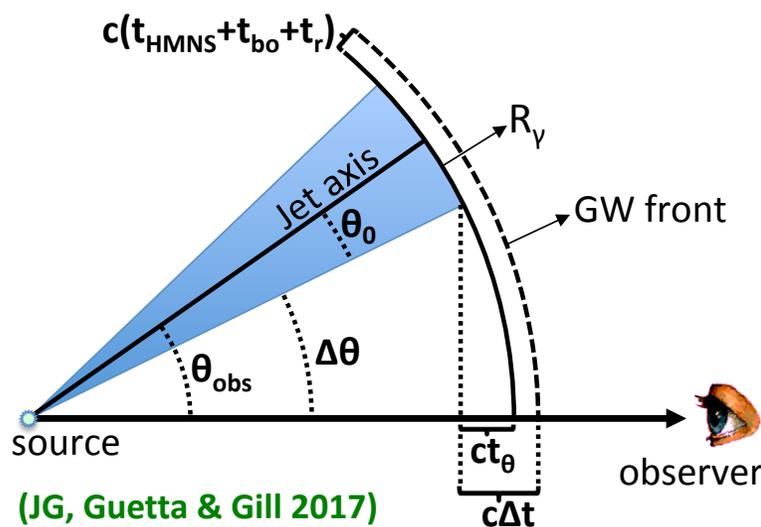
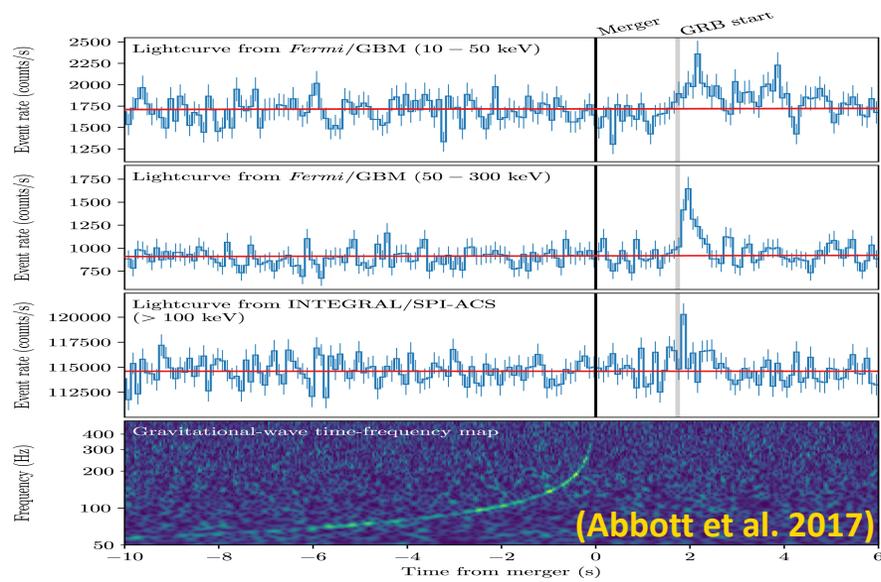
GW170817 / GRB170817A: The Time Delay

- The $\Delta t \approx 1.74$ s delay between the GW chirp signal & the sGRB onset $\Rightarrow \left| \frac{v_{GW}}{c} - 1 \right| \lesssim 4 \cdot 10^{-16}$
- A HMNS may explain $\Delta t \approx 1.74$ s by $t_{HMNS} \lesssim 0.5$ s & $t_{bo} \sim 1$ s
(Moharana & Piran 2017 find $t_{bo} \sim 0.5$ s for SGRBs, from a plateau in their duration distribution, dN_{GRB}/dT_{GRB})
- Direct BH formation \Rightarrow a shorter jet breakout time $t_{bo} \Rightarrow$ the jet is less likely to be choked
- If the prompt γ -rays are beamed away from us (large $\Gamma\Delta\theta$), the implied on-axis $L_{\gamma,iso}$ & E_{peak} are very high – inconsistent with their observed correlation (JG+ 2017) & implying large compactness (Matsumoto+ 2019) \Rightarrow they must arise from $\Gamma\Delta\theta < 1 \Rightarrow$ a jet with angular structure



GW170817 / GRB170817A: The Time Delay

- The $\Delta t \approx 1.74$ s delay between the GW chirp signal & the sGRB onset $\Rightarrow \left| \frac{v_{GW}}{c} - 1 \right| \lesssim 4 \cdot 10^{-16}$
- A HMNS may explain $\Delta t \approx 1.74$ s by $t_{HMNS} \lesssim 0.5$ s & $t_{bo} \sim 1$ s
(Moharana & Piran 2017 find $t_{bo} \sim 0.5$ s for SGRBs, from a plateau in their duration distribution, dN_{GRB}/dT_{GRB})
- Direct BH formation \Rightarrow a shorter jet breakout time $t_{bo} \Rightarrow$ the jet is less likely to be choked
- If the prompt γ -rays are beamed away from us (large $\Gamma\Delta\theta$), the implied on-axis $L_{\gamma,iso}$ & E_{peak} are very high – inconsistent with their observed correlation (JG+ 2017) & implying large compactness (Matsumoto+ 2019) \Rightarrow they must arise from $\Gamma\Delta\theta < 1 \Rightarrow$ a jet with angular structure



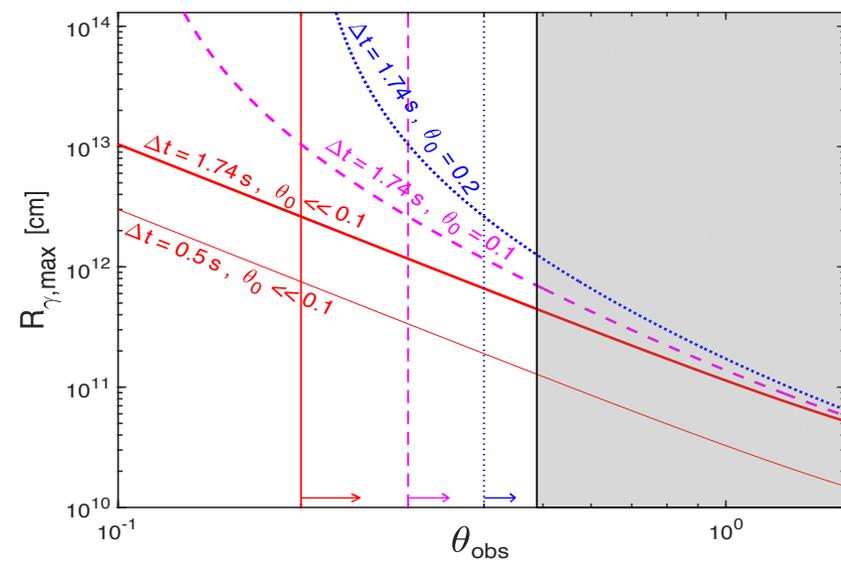
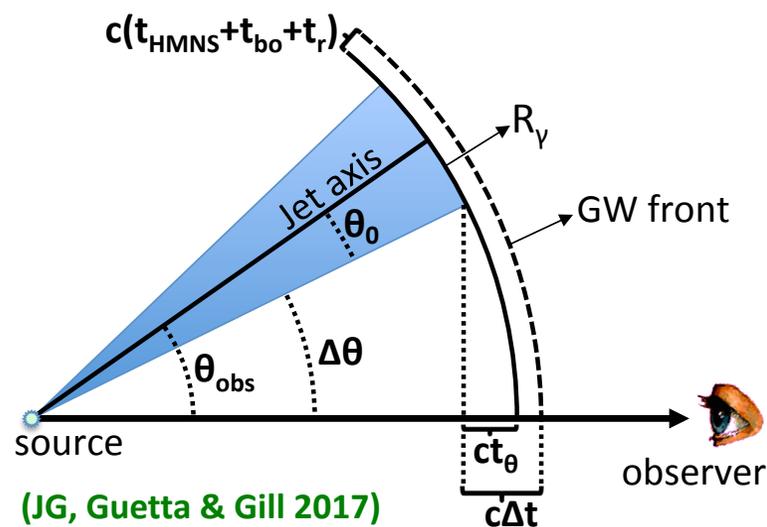
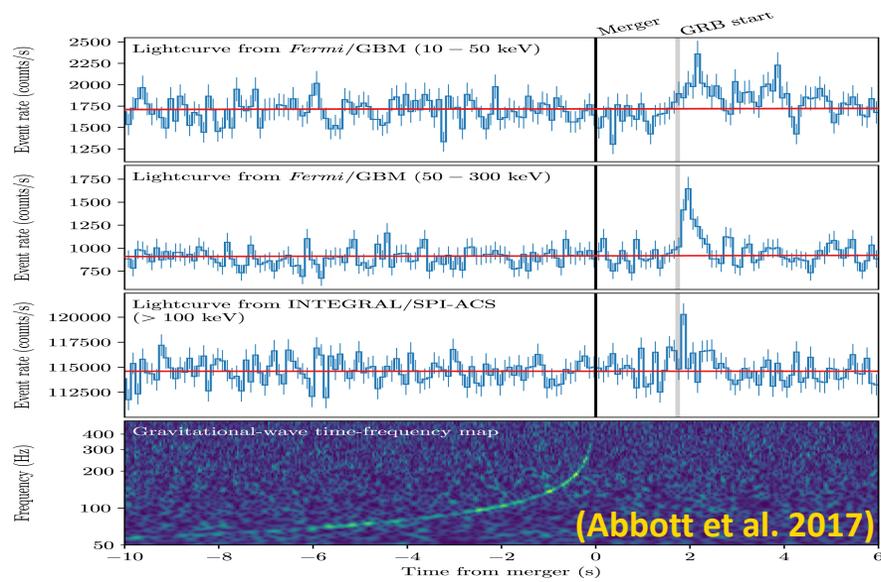
$$t_{\theta} = \frac{R_{\gamma}}{c} [1 - \cos(\Delta\theta)] \approx \frac{R_{\gamma}}{2c} \Delta\theta^2$$

$$R_{\gamma} < \frac{c\Delta t}{1 - \cos(\Delta\theta)} \approx \frac{2c\Delta t}{\Delta\theta^2}$$

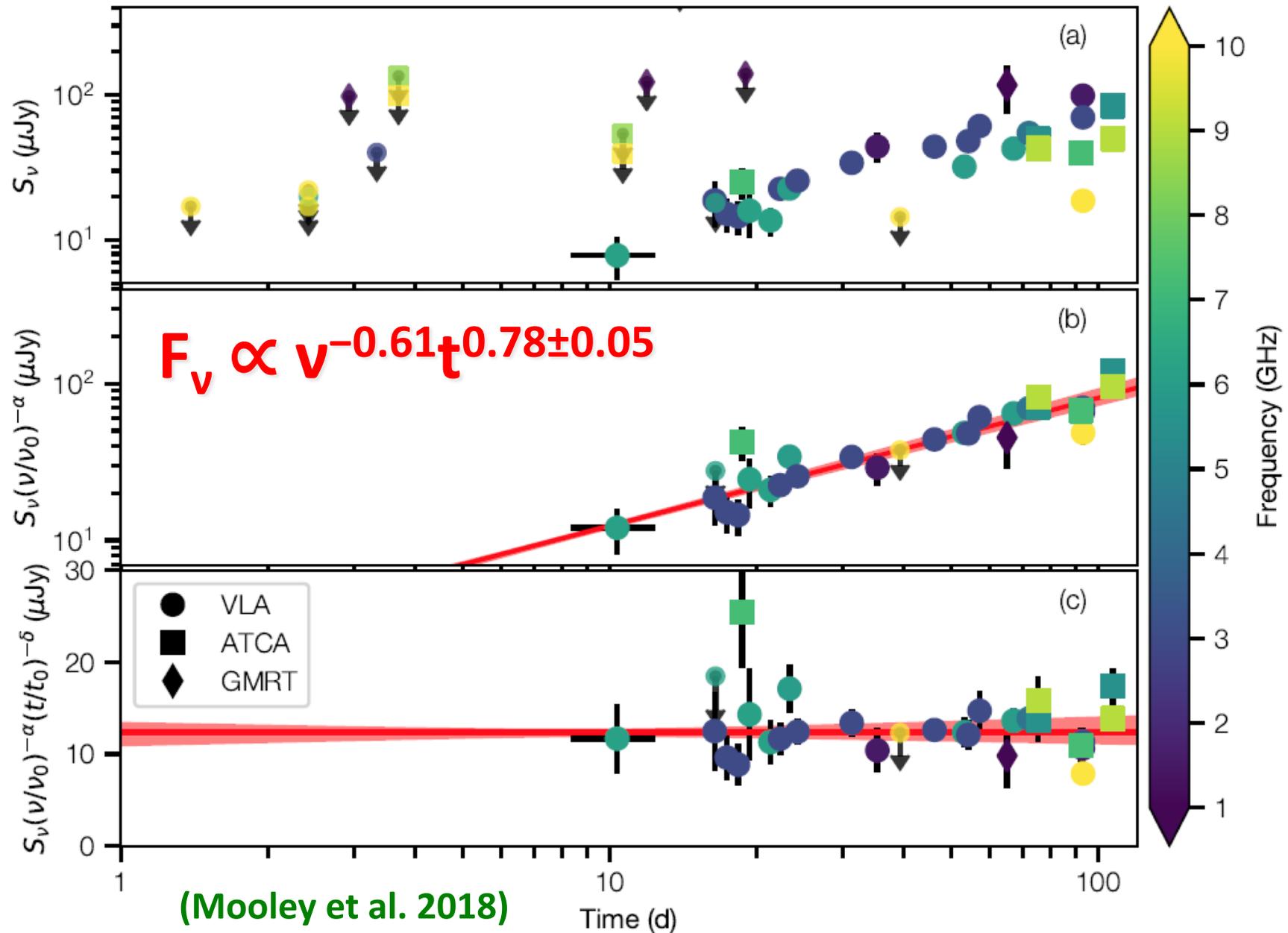
$$t_r = t_{lab} - \frac{R_{\gamma}}{c} \approx \frac{R_{\gamma}}{2c\Gamma^2}$$

GW170817 / GRB170817A: The Time Delay

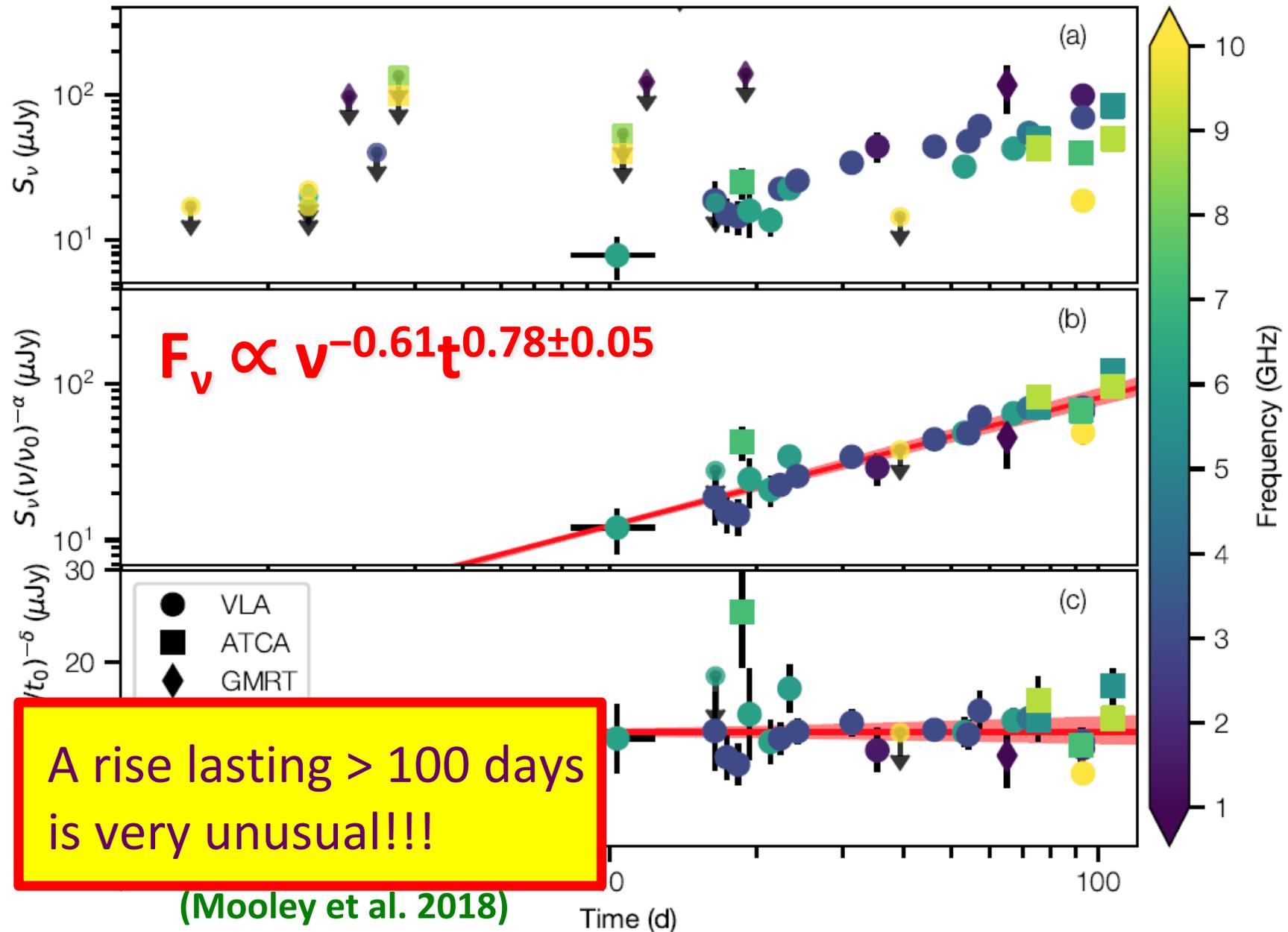
- The $\Delta t \approx 1.74$ s delay between the GW chirp signal & the sGRB onset $\Rightarrow \left| \frac{v_{GW}}{c} - 1 \right| \lesssim 4 \cdot 10^{-16}$
- A HMNS may explain $\Delta t \approx 1.74$ s by $t_{HMNS} \lesssim 0.5$ s & $t_{bo} \sim 1$ s
(Moharana & Piran 2017 find $t_{bo} \sim 0.5$ s for SGRBs, from a plateau in their duration distribution, dN_{GRB}/dT_{GRB})
- Direct BH formation \Rightarrow a shorter jet breakout time $t_{bo} \Rightarrow$ the jet is less likely to be choked
- If the prompt γ -rays are beamed away from us (large $\Gamma\Delta\theta$), the implied on-axis $L_{\gamma,iso}$ & E_{peak} are very high – inconsistent with their observed correlation (JG+ 2017) & implying large compactness (Matsumoto+ 2019) \Rightarrow they must arise from $\Gamma\Delta\theta < 1 \Rightarrow$ a jet with angular structure



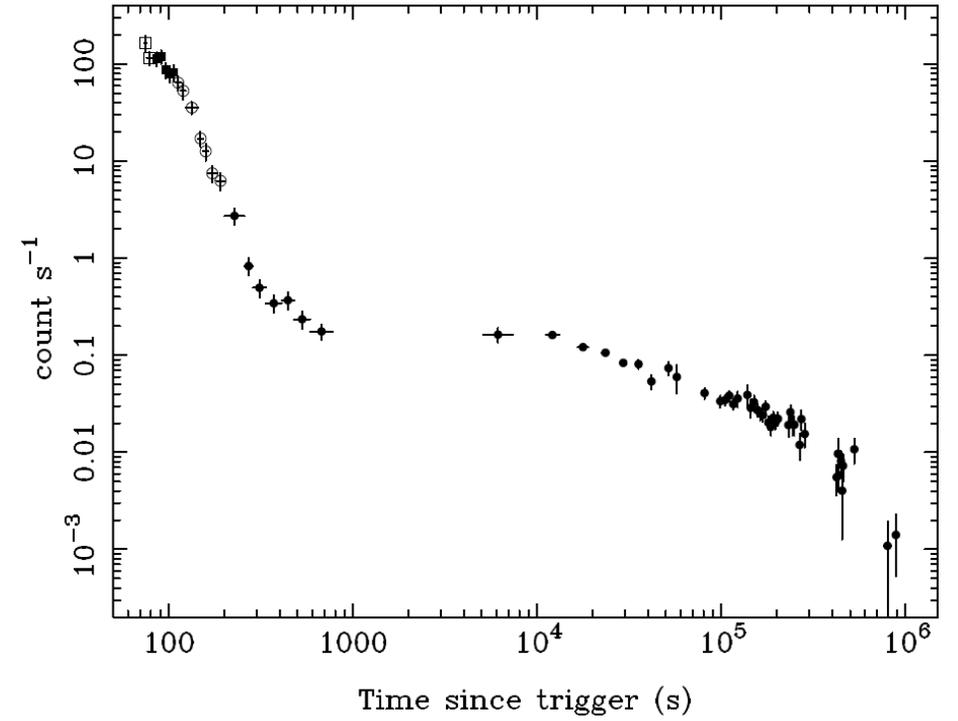
GRB170817A: Afterglow Observations



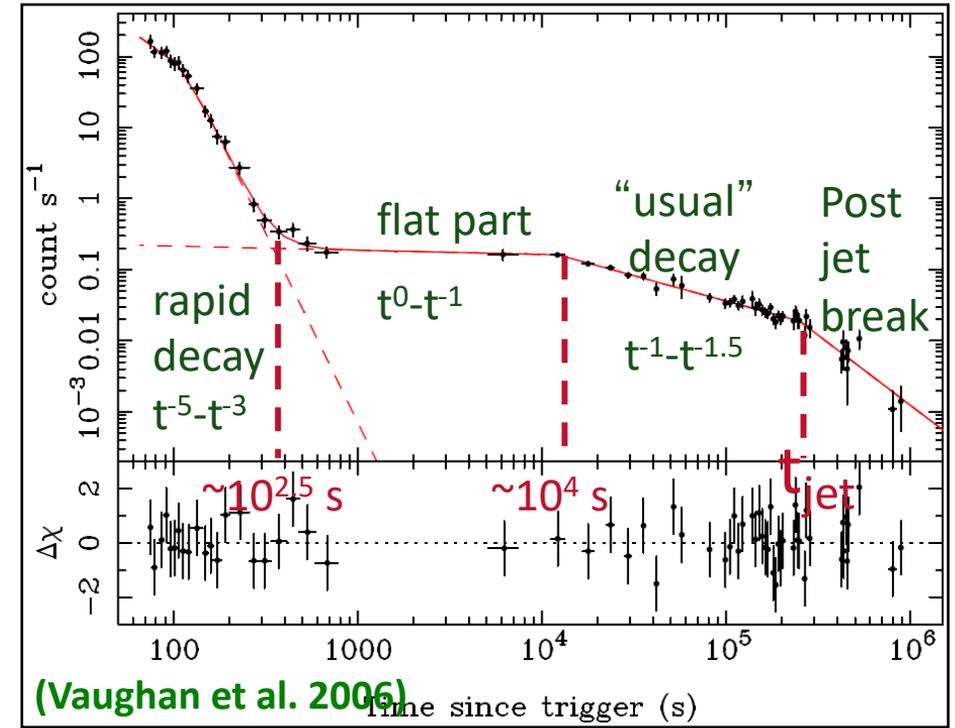
GRB170817A: Afterglow Observations



Analogy to rising F_ν : X-ray Plateaus

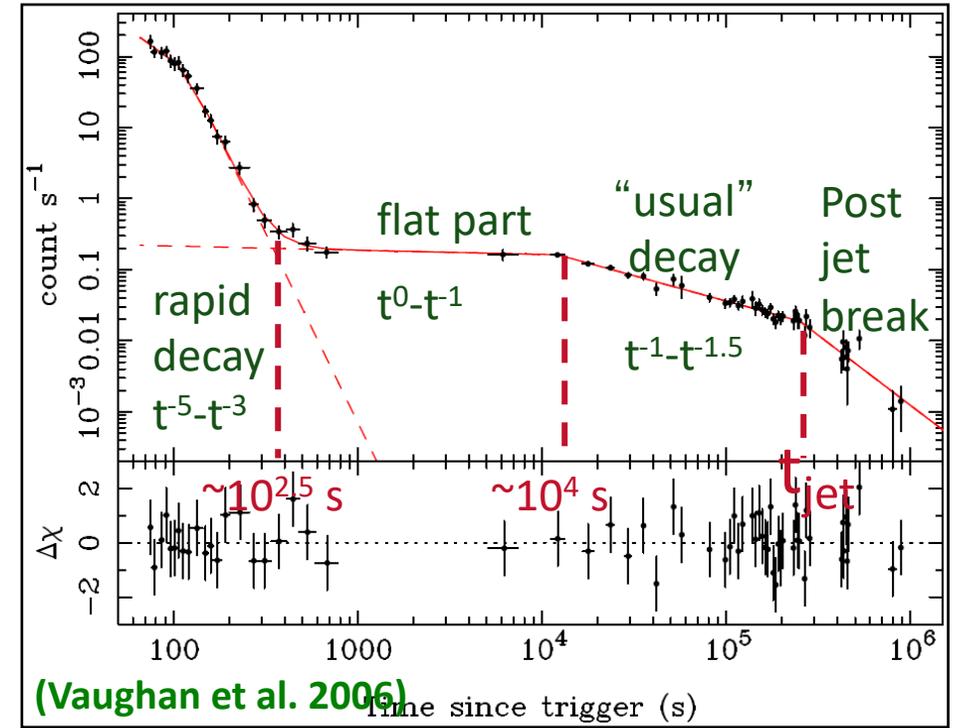


Analogy to rising F_ν : X-ray Plateaus



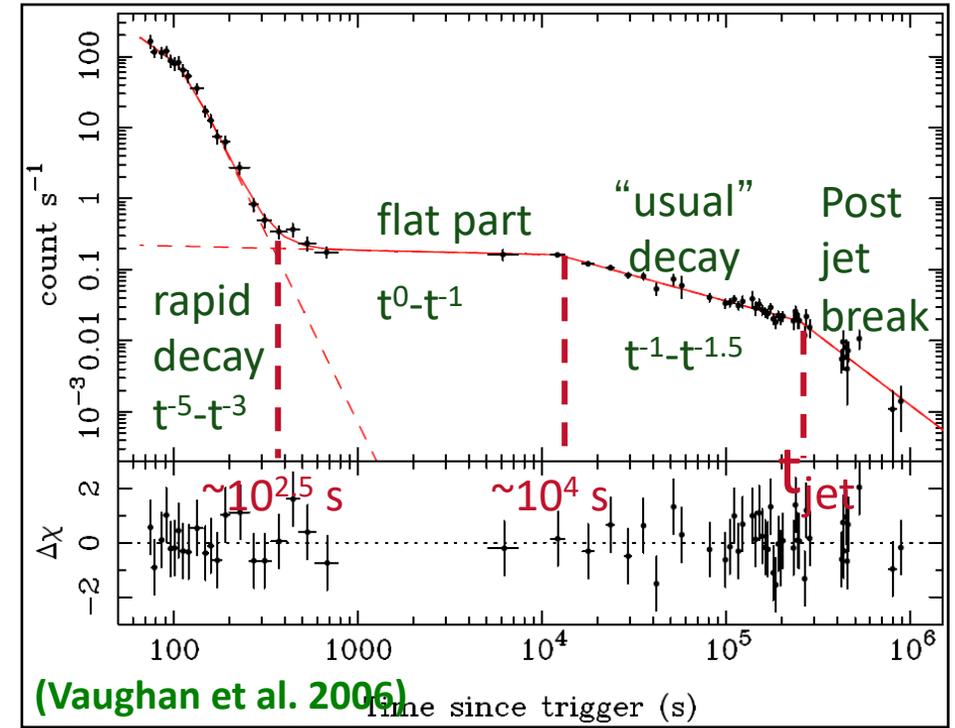
Analogy to rising F_ν : X-ray Plateaus

- Possible solutions:
- ◆ Evolution of shock microphysical parameters (JG, Konigl & Piran 2006)



Analogy to rising F_ν : X-ray Plateaus

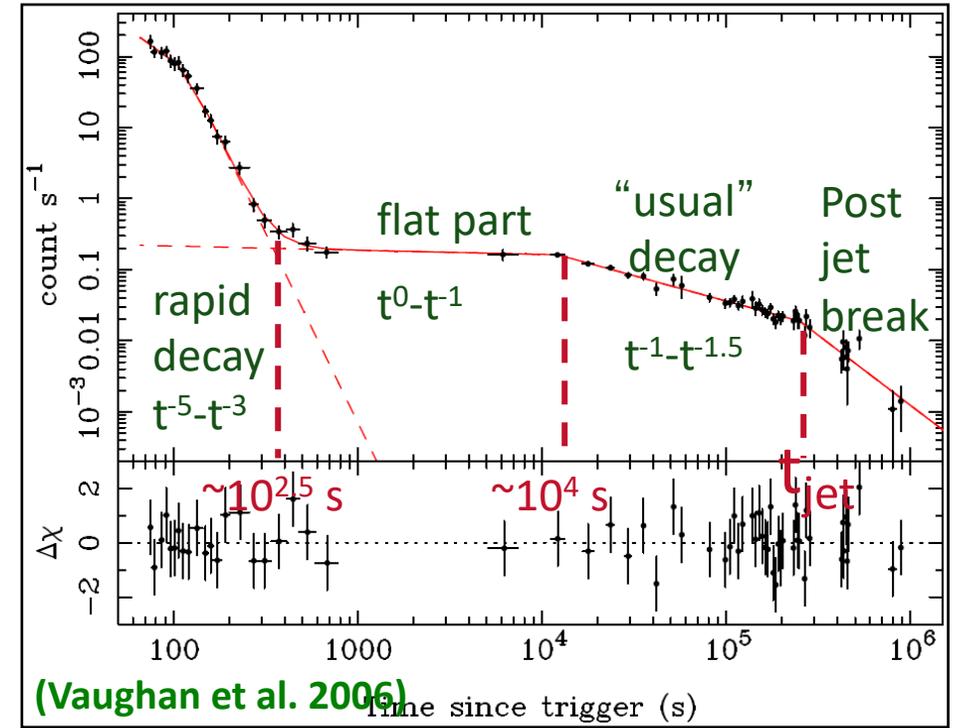
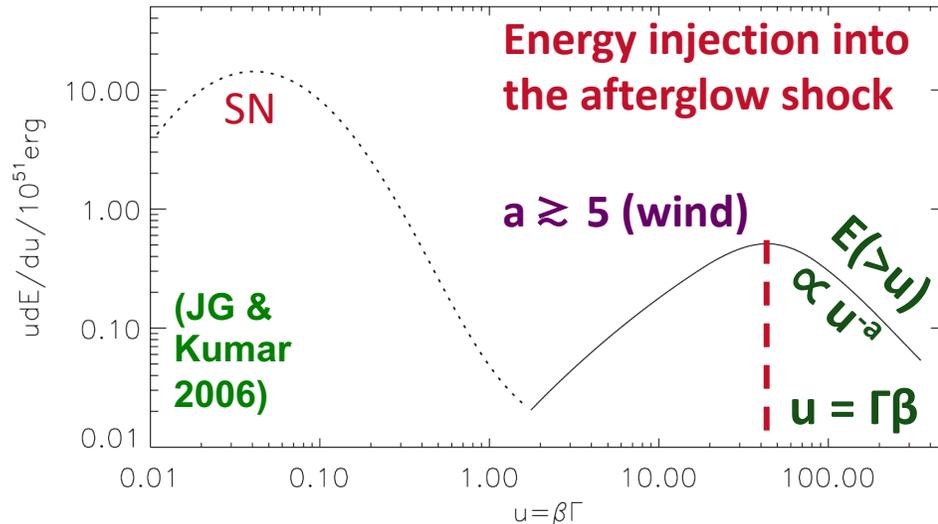
- Possible solutions:
 - ◆ Evolution of shock microphysical parameters (JG, Konigl & Piran 2006)
 - ◆ Energy injection into external shock:
 1. long-lived relativistic wind



Analogy to rising F_ν : X-ray Plateaus

■ Possible solutions:

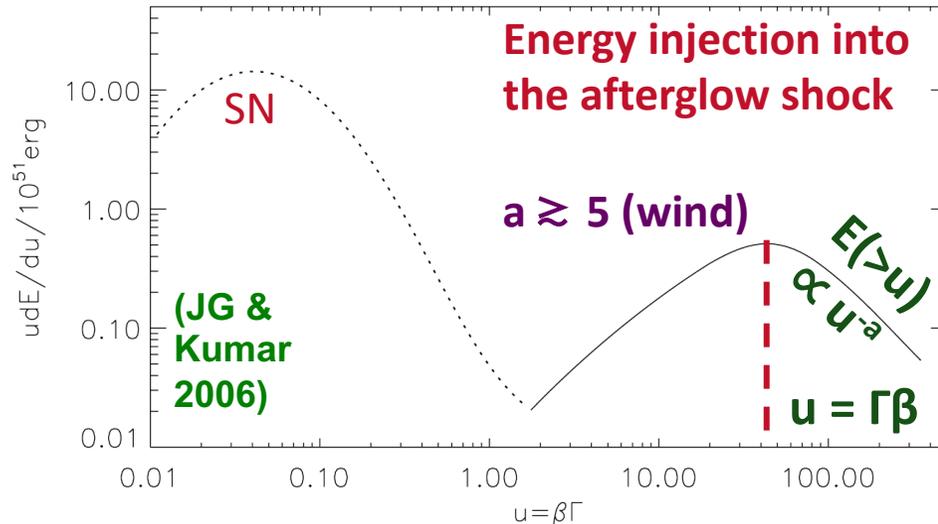
- ◆ Evolution of shock microphysical parameters (JG, Konigl & Piran 2006)
- ◆ Energy injection into external shock:
 1. long-lived relativistic wind
 2. slower ejecta catching up
(Sari & Meszaros 00; Nousek+ 06; JG & Kumar 06)



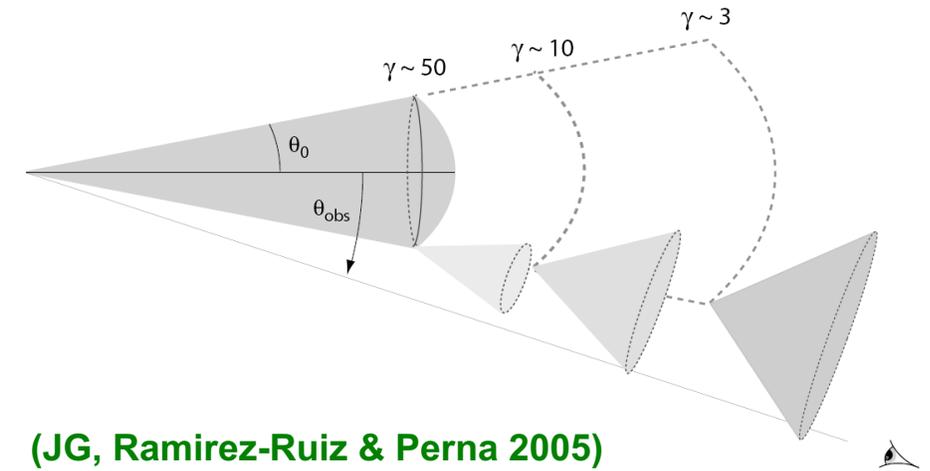
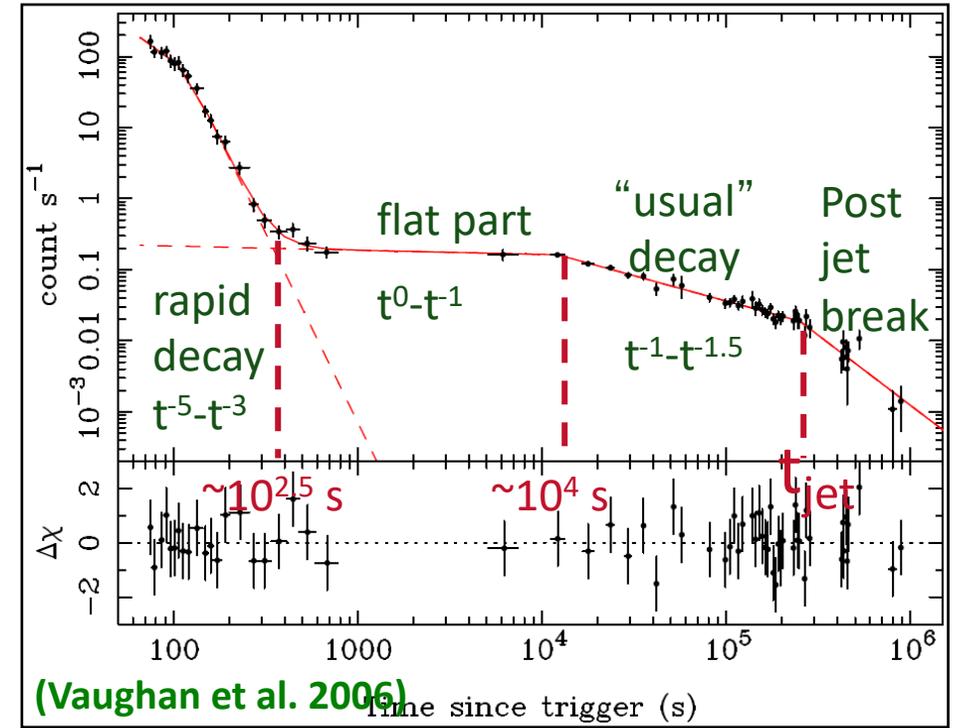
Analogy to rising F_ν : X-ray Plateaus

■ Possible solutions:

- ◆ Evolution of shock microphysical parameters (JG, Konigl & Piran 2006)
- ◆ Energy injection into external shock:
 1. long-lived relativistic wind
 2. slower ejecta catching up
(Sari & Meszaros 00; Nousek+ 06; JG & Kumar 06)



◆ Viewing angle effects

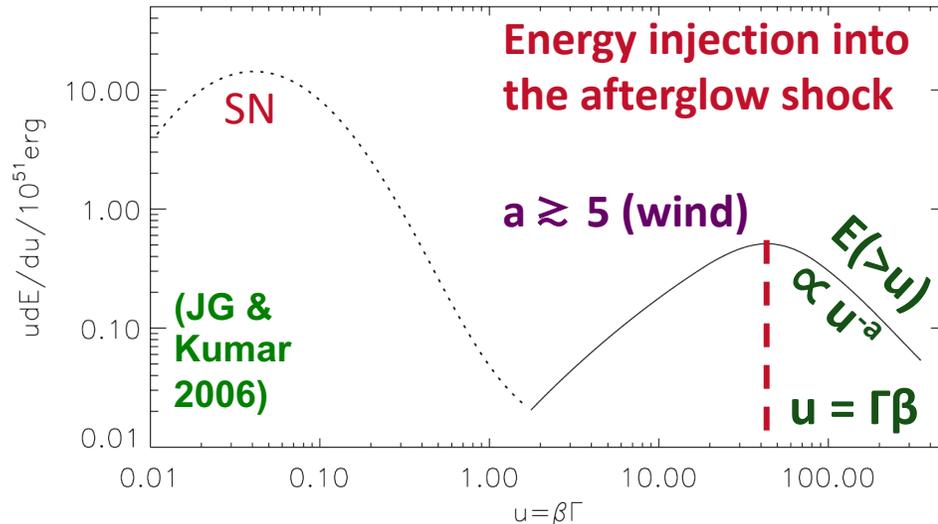


(JG, Ramirez-Ruiz & Perna 2005)

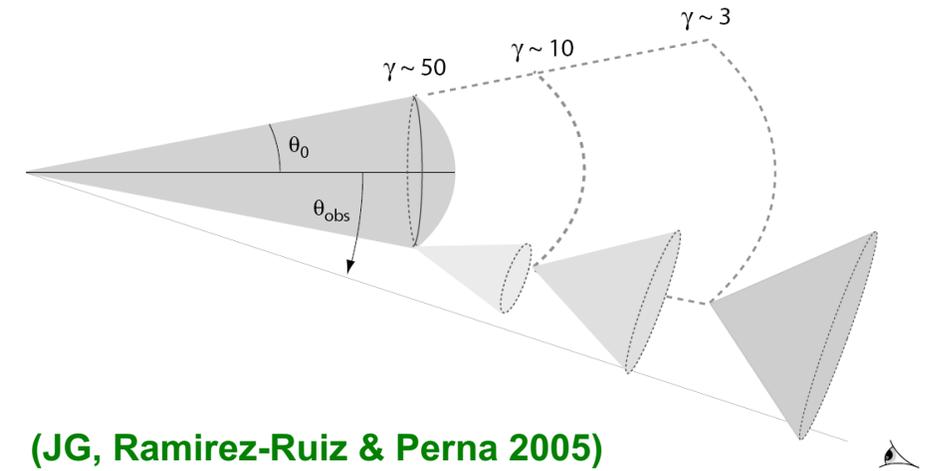
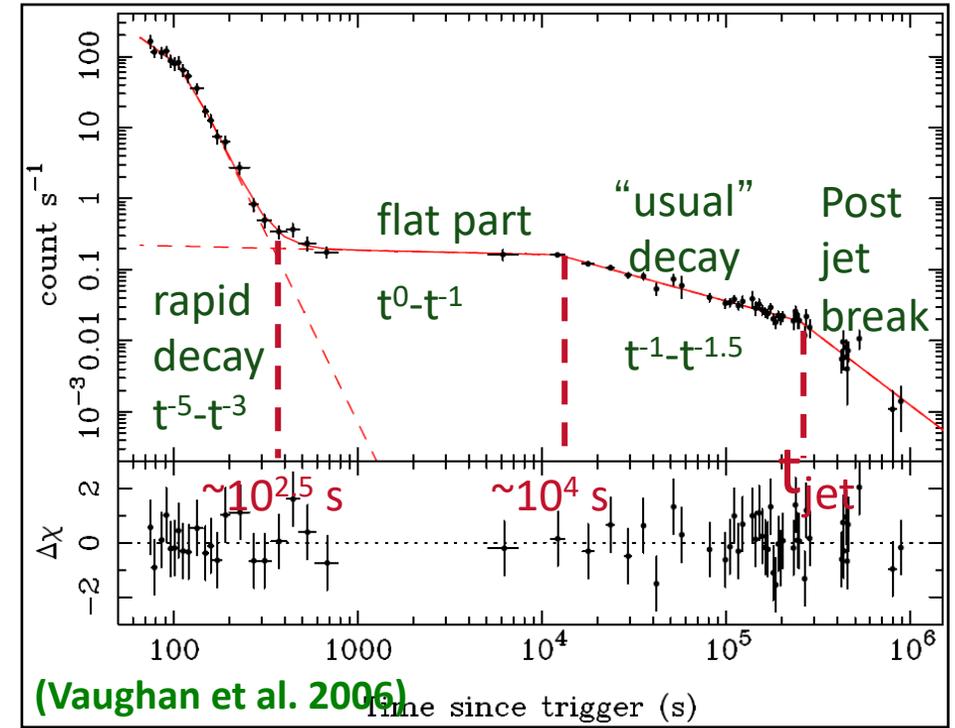
Analogy to rising F_ν : X-ray Plateaus

■ Possible solutions:

- ◆ Evolution of shock microphysical parameters (JG, Konigl & Piran 2006)
- ◆ Energy injection into external shock:
 1. long-lived relativistic wind
 2. **slower ejecta catching up** radial
(Sari & Meszaros 00; Nousek+ 06; JG & Kumar 06)



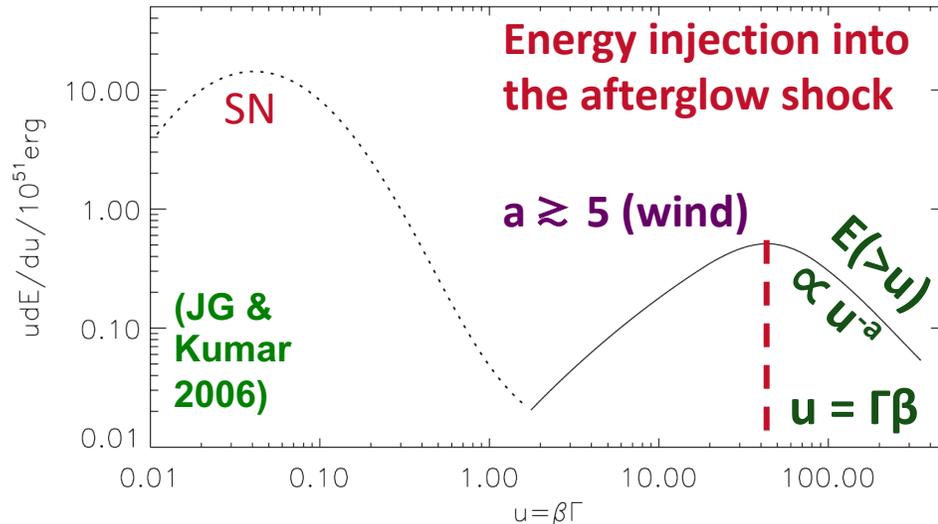
◆ Viewing angle effects



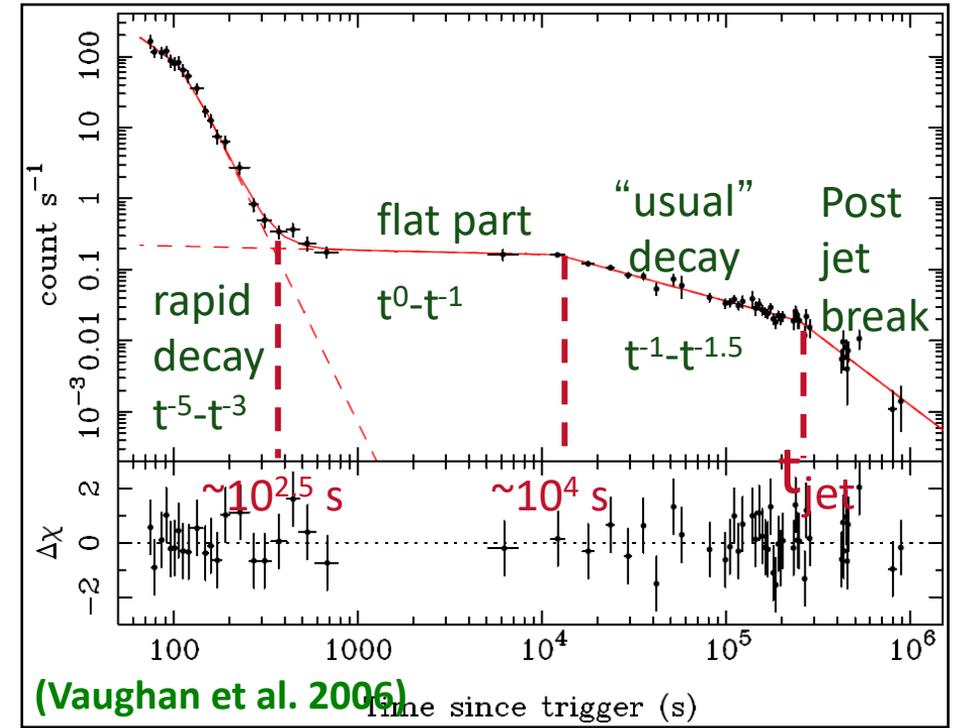
Analogy to rising F_ν : X-ray Plateaus

■ Possible solutions:

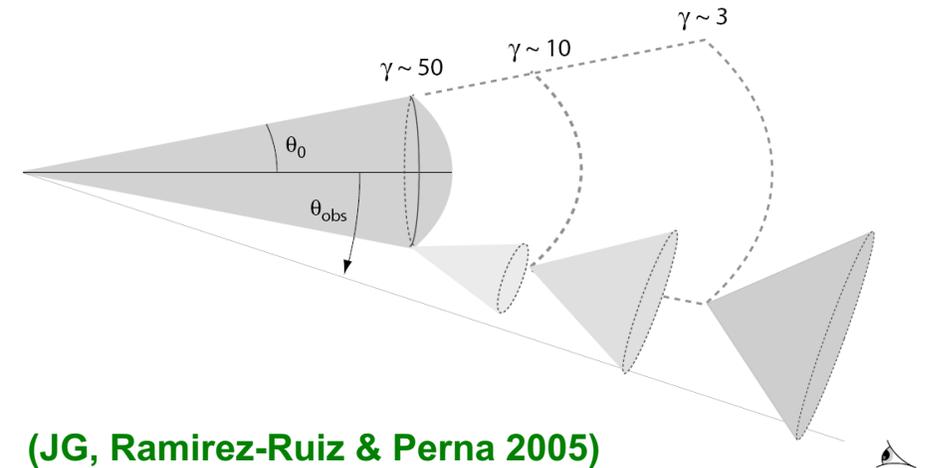
- ◆ Evolution of shock microphysical parameters (JG, Konigl & Piran 2006)
- ◆ Energy injection into external shock:
 1. long-lived relativistic wind
 2. **slower ejecta catching up** radial
(Sari & Meszaros 00; Nousek+ 06; JG & Kumar 06)



- ◆ **Viewing angle effects** angular



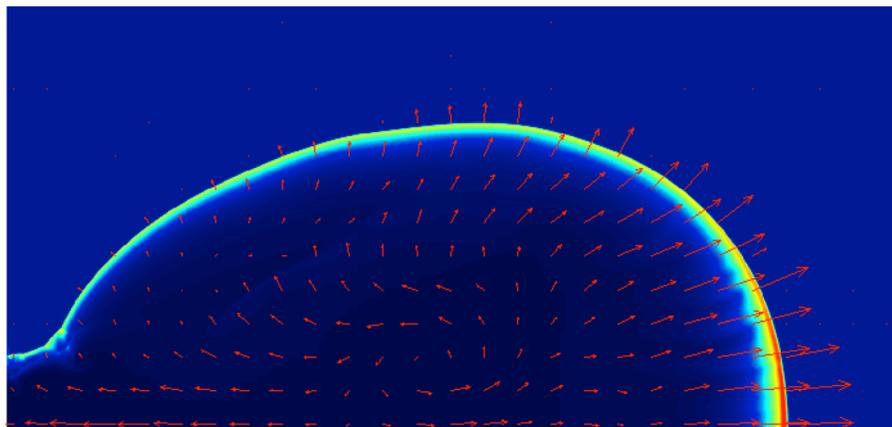
(Vaughan et al. 2006)



(JG, Ramirez-Ruiz & Perna 2005)

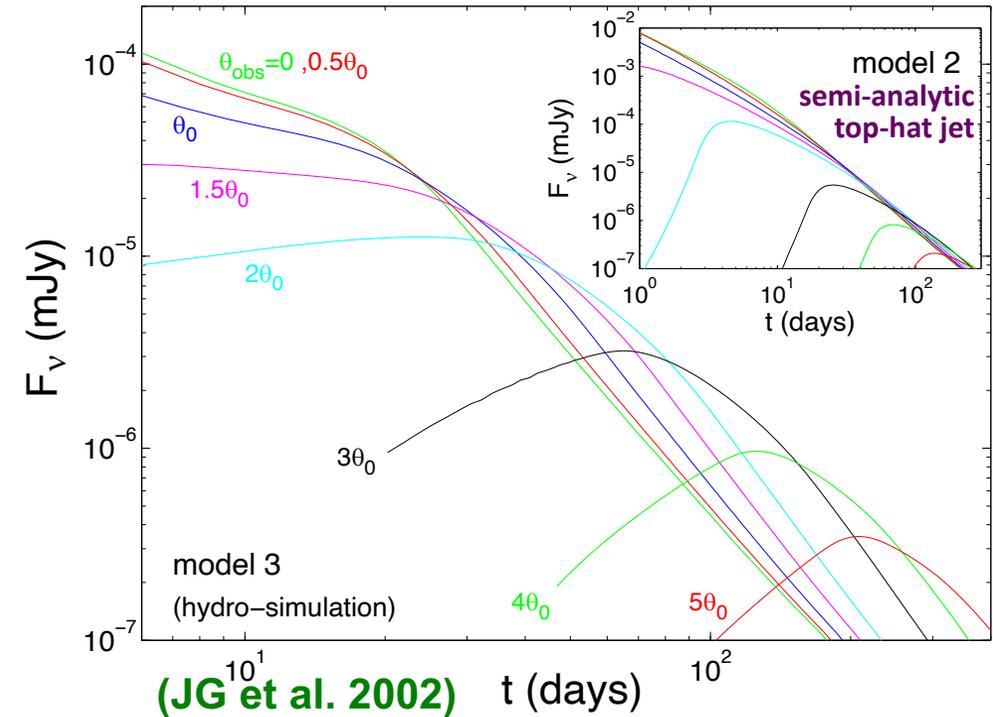
Analogy to rising F_{ν} : Off-Axis Viewing

- The emission is initially strongly beamed away from our L.o.S
- F_{ν} rises as beaming cone widens
- When beaming cone reaches LoS F_{ν} peaks & approaches on-axis F_{ν}
- The rise is much more gradual for hydrodynamic simulations due to slower matter at the jet's sides with non-radial velocities

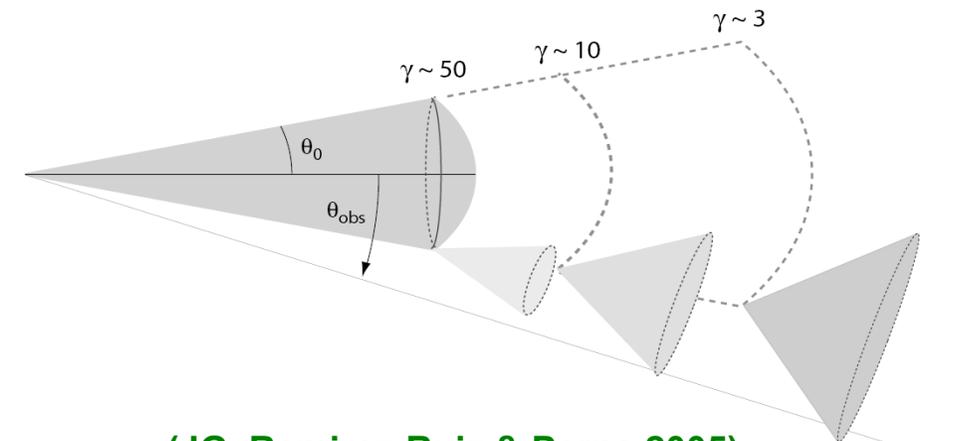


(JG et al. 2001)

0.311E-12 0.163E-06



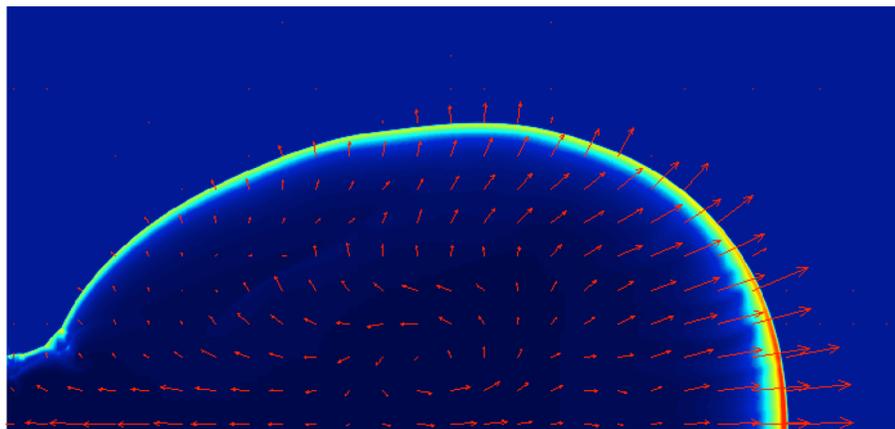
(JG et al. 2002)



(JG, Ramirez-Ruiz & Perna 2005)

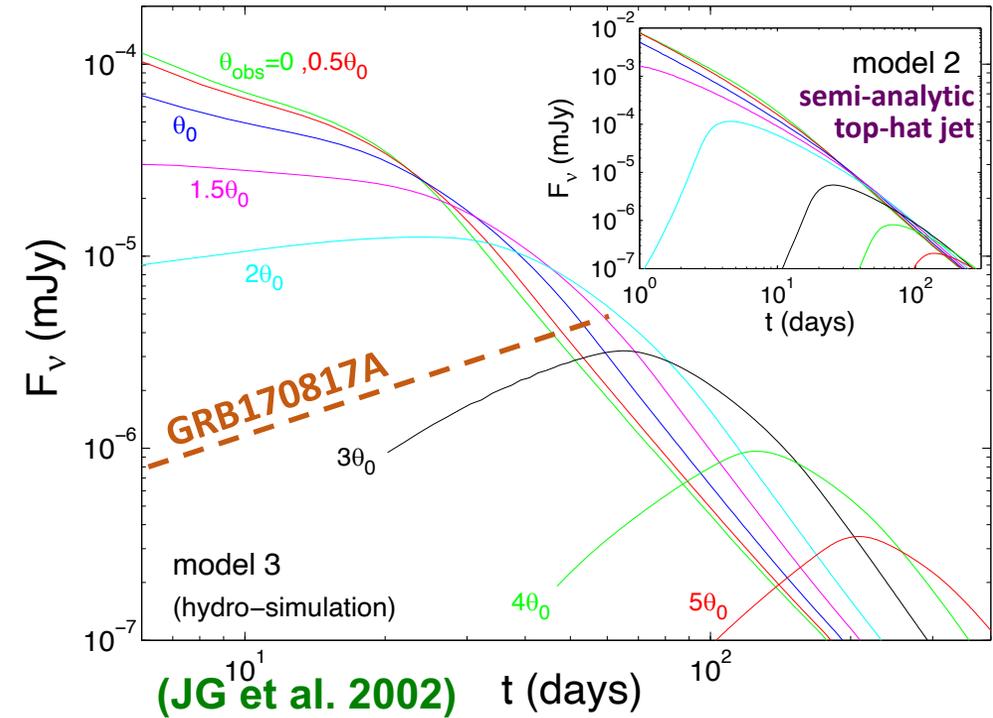
Analogy to rising F_ν : Off-Axis Viewing

- The emission is initially strongly beamed away from our L.o.S
- F_ν rises as beaming cone widens
- When beaming cone reaches LoS F_ν peaks & approaches on-axis F_ν
- The rise is much more gradual for hydrodynamic simulations due to slower matter at the jet's sides with non-radial velocities

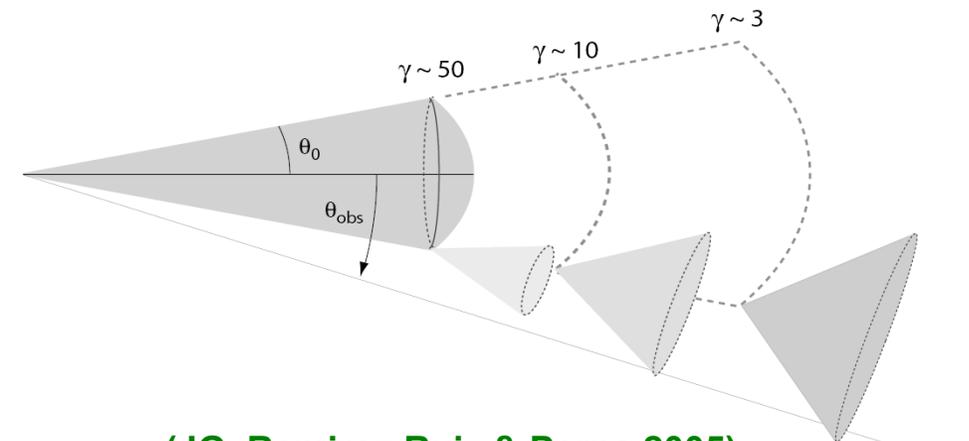


(JG et al. 2001)

0.311E-12 0.163E-06



(JG et al. 2002)

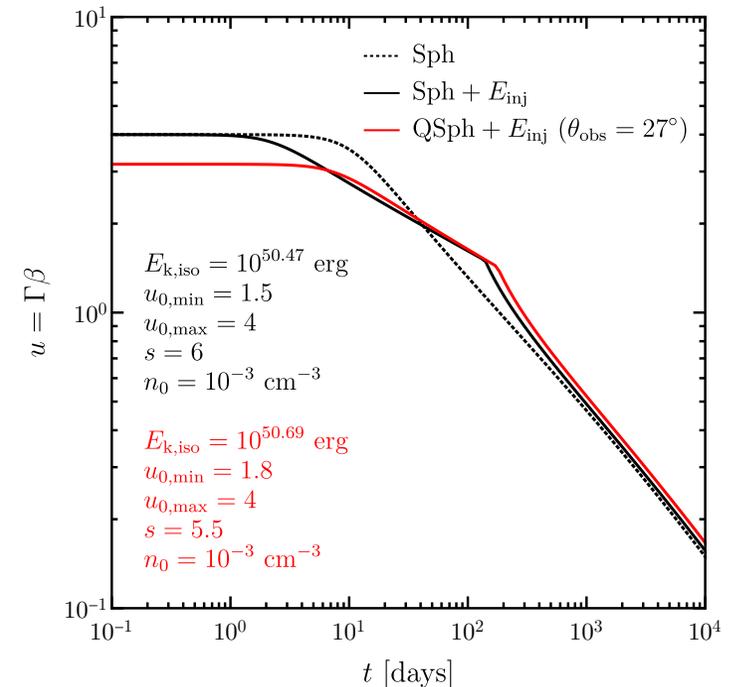
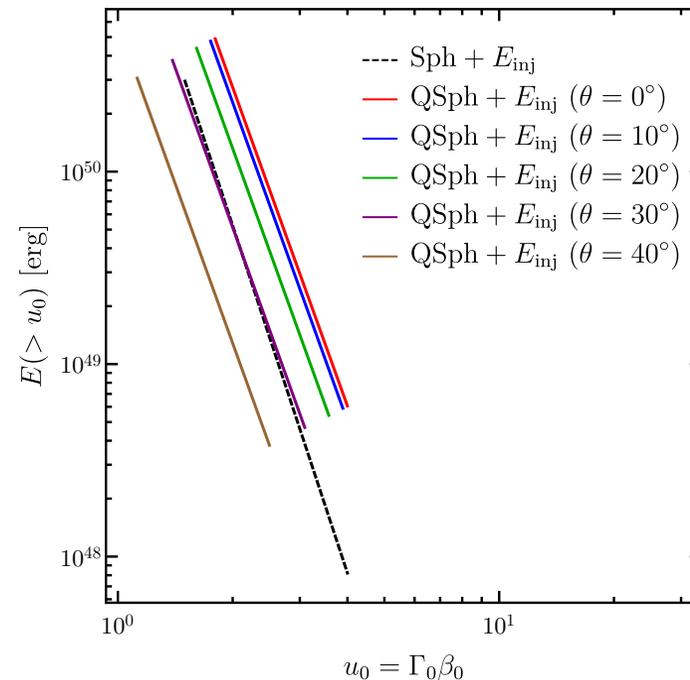


(JG, Ramirez-Ruiz & Perna 2005)

Outflow Structure: Breaking the Degeneracy (JG & Gill 18)

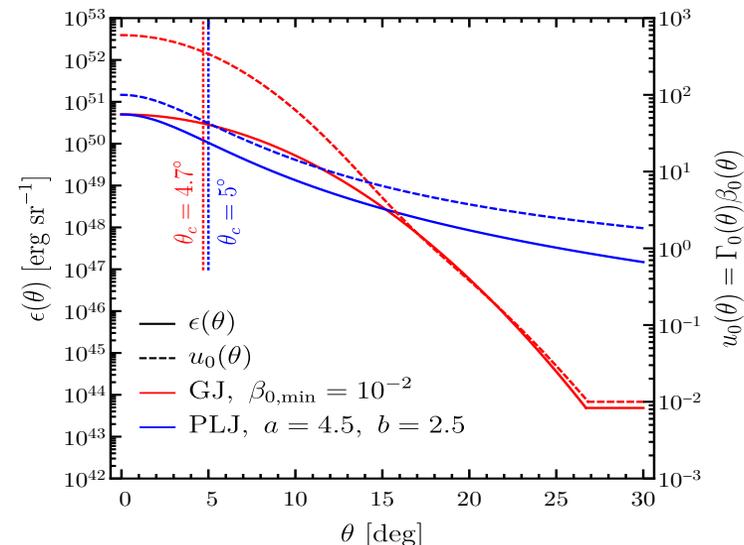
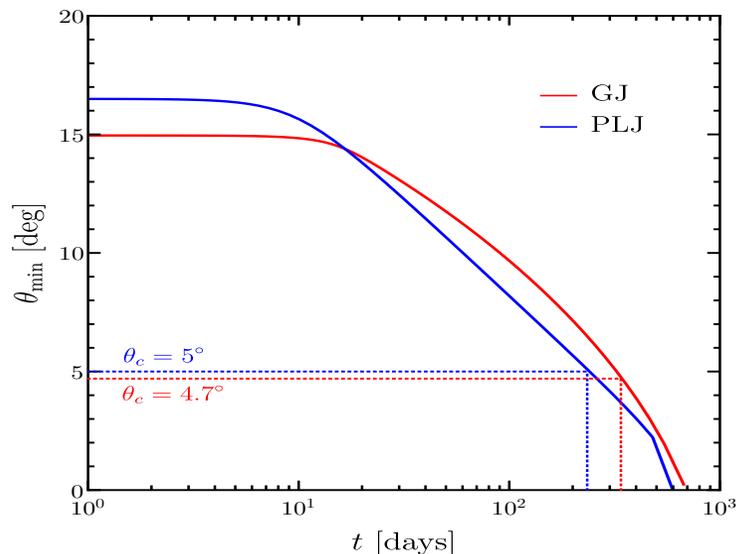
- The lightcurves leave a lot of degeneracy between models
- The degeneracy may be lifted by calculation the afterglow **images** & **polarization** (e.g. Nakar & Piran 2018; Nakar et al. 2018)
- We considered 4 different models including both main types
 - ◆ Sph+ E_{inj} : Spherical with energy injection $E(>u=\Gamma\beta) \propto u^{-6}$, $1.5 < u < 4$
 - ◆ QSph+ E_{inj} : Quasi-Spherical + energy injection $E(>u) \propto u^{-s}$, $u_{min,0} = 1.8$ $u_{max,0} = 4$, $s = 5.5$, $\zeta = 0.1$

$$\frac{\epsilon(\theta)}{\epsilon_0} = \frac{u_{0,min}(\theta)}{u_{min,0}} = \frac{u_{0,max}(\theta)}{u_{max,0}} = \frac{\zeta + \cos^2 \theta}{\zeta + 1}$$



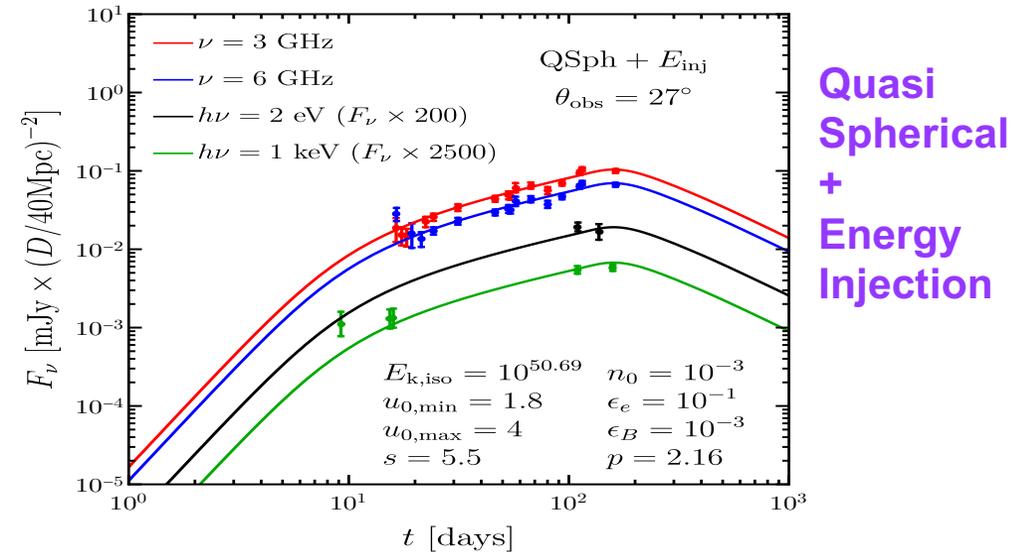
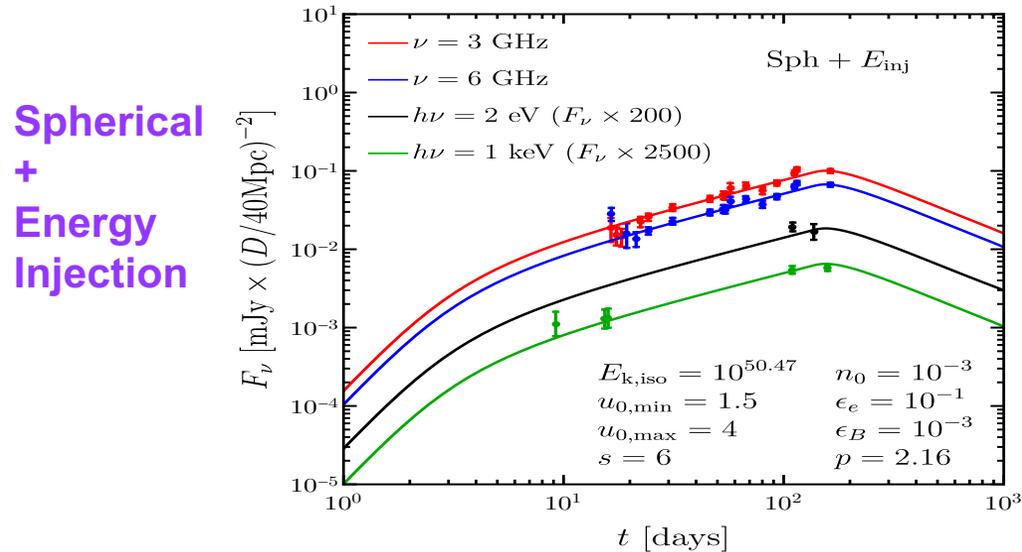
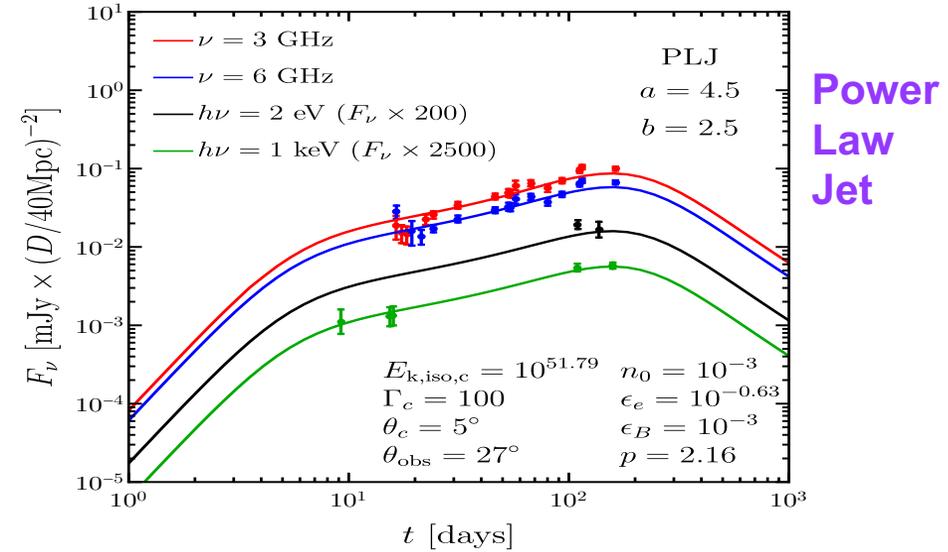
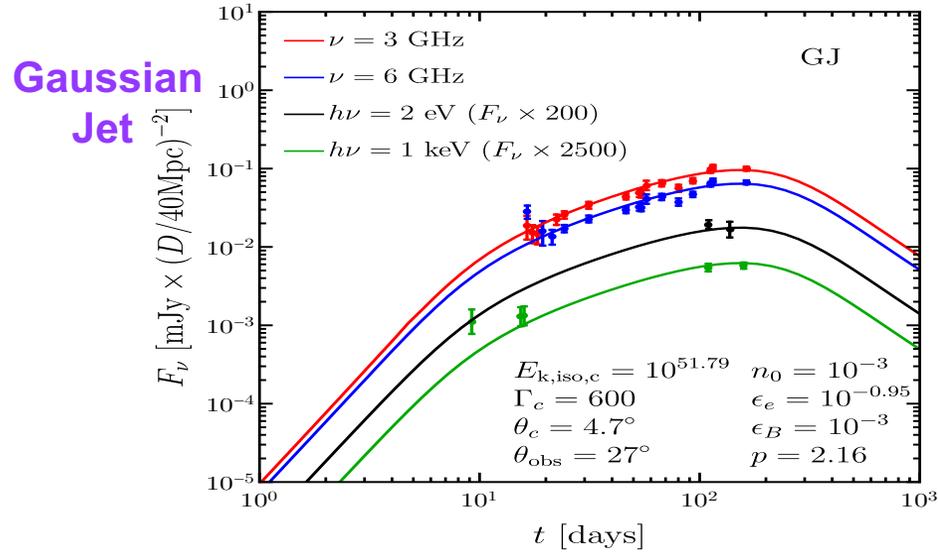
Outflow Structure: Breaking the Degeneracy (JG & Gill 18)

- The lightcurves leave a lot of degeneracy between models
- The degeneracy may be lifted by calculation the afterglow **images** & **polarization** (e.g. Nakar & Piran 2018; Nakar et al. 2018)
- We considered 4 different models including both main types
 - ◆ GJ: Gaussian Jet (in $\epsilon = dE/d\Omega$, $\Gamma_0 - 1$) $\Gamma_c = 600$, $\theta_c = 4.7^\circ$
 - ◆ PLJ: Power-Law Jet; $\epsilon = \epsilon_c \Theta^{-a}$, $\Gamma_0 - 1 = (\Gamma_c - 1) \Theta^{-b}$, $\Theta = [1 + (\theta/\theta_c)^2]^{1/2}$, $\Gamma_c = 100$, $\theta_c = 5^\circ$, $a = 4.5$, $b = 2.5$
- As there is a lot of freedom we fixed: $p = 2.16$, $\epsilon_B = n_0 = 10^{-3}$, $\theta_{\text{obs}} = 27^\circ$



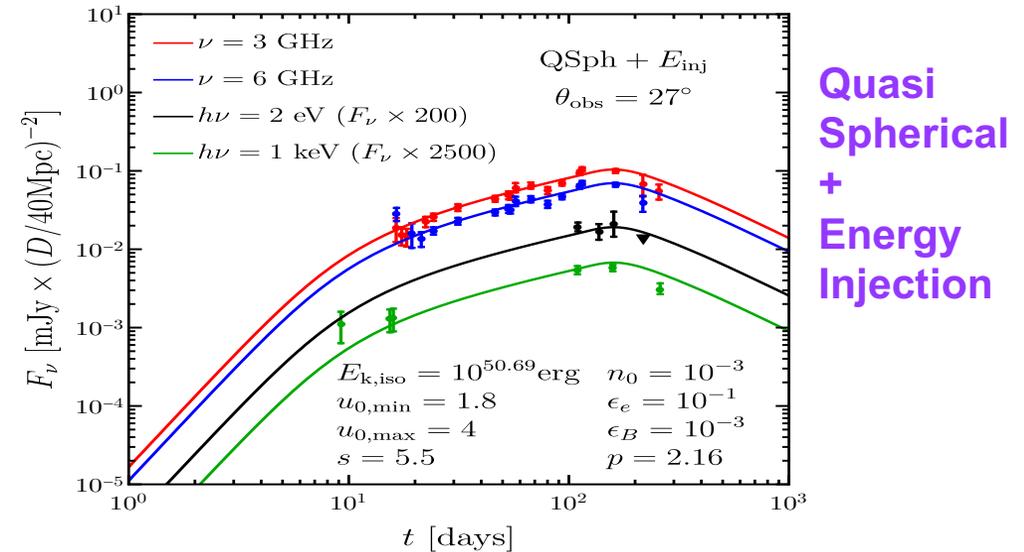
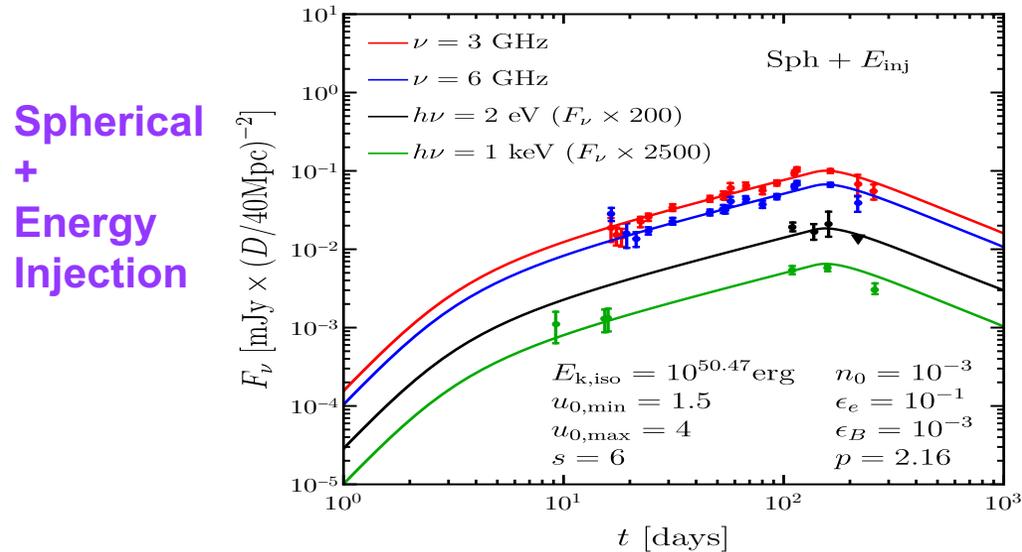
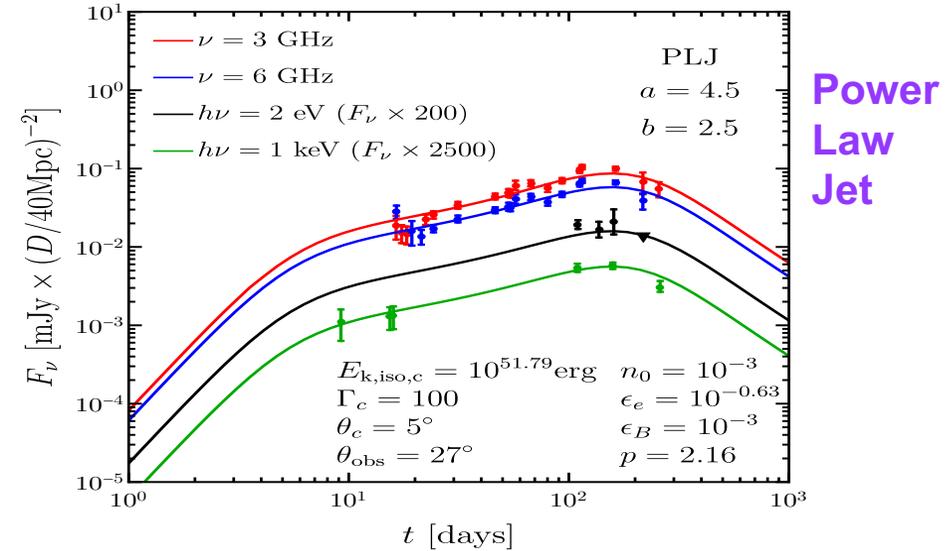
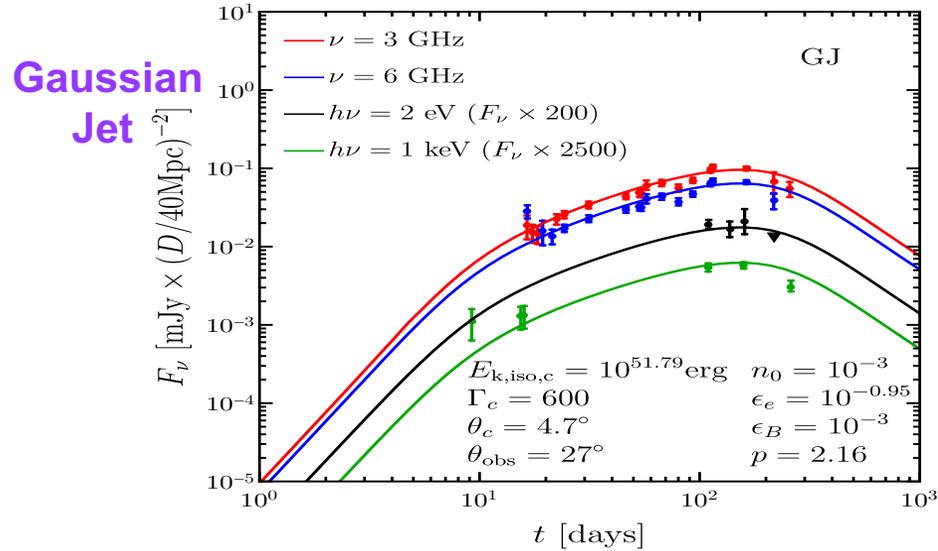
Outflow Structure: Breaking the Degeneracy (JG & Gill 18)

■ Tentative fit to GRB170817A afterglow data (radio to X-ray)



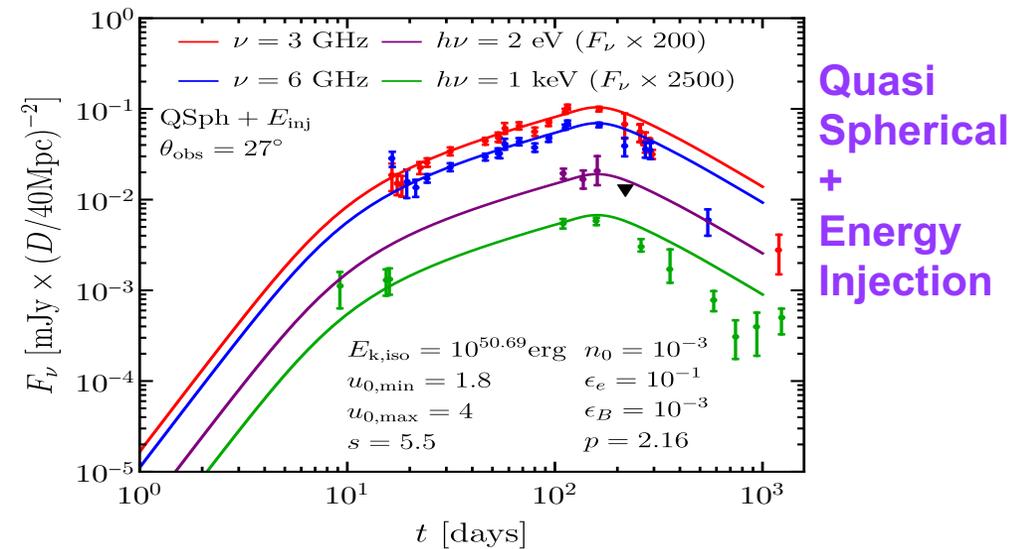
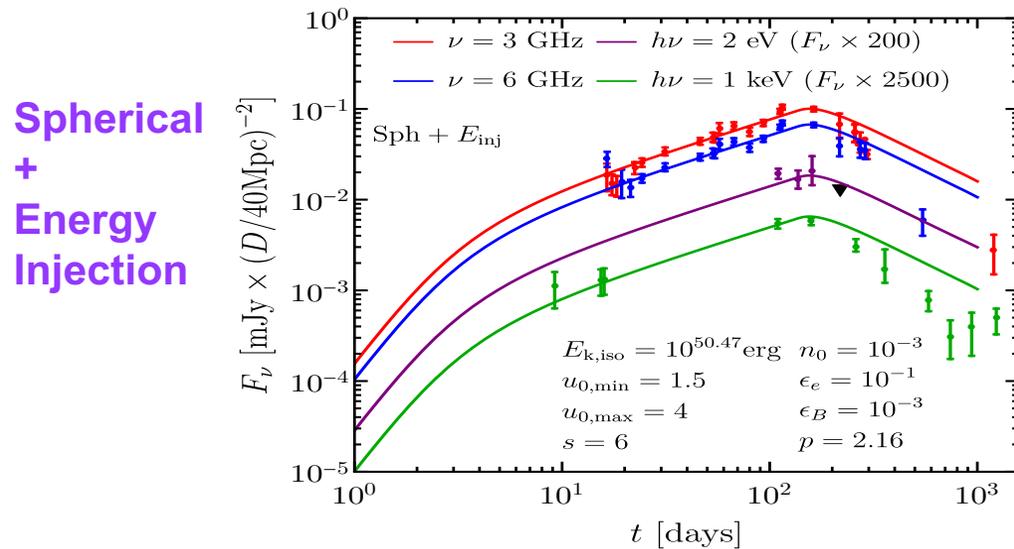
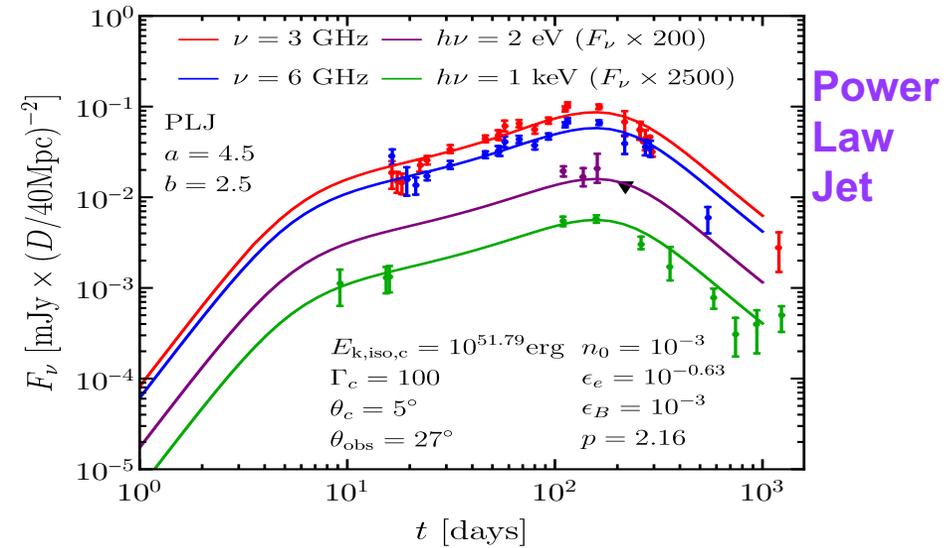
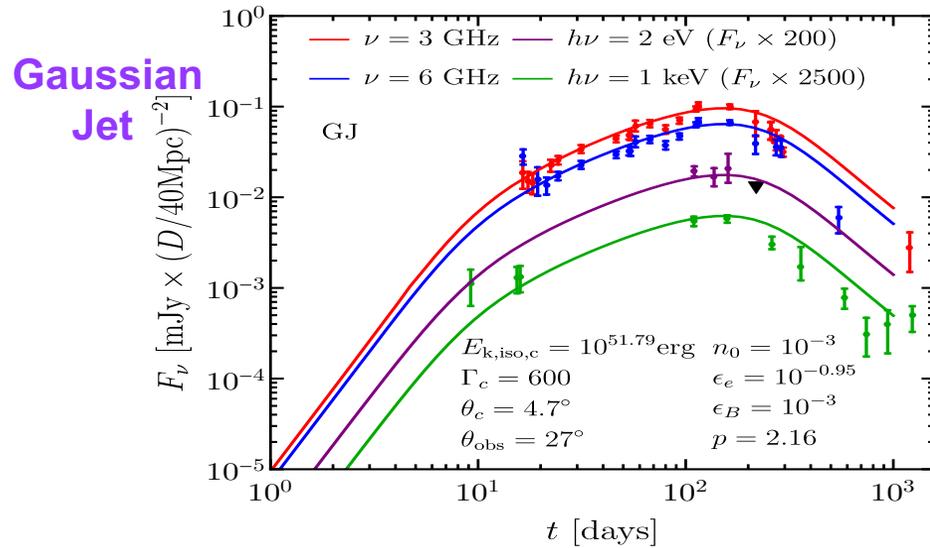
Outflow Structure: Breaking the Degeneracy (JG & Gill 18)

- New data that came out established a peak at $t_{\text{peak}} \sim 150$ days



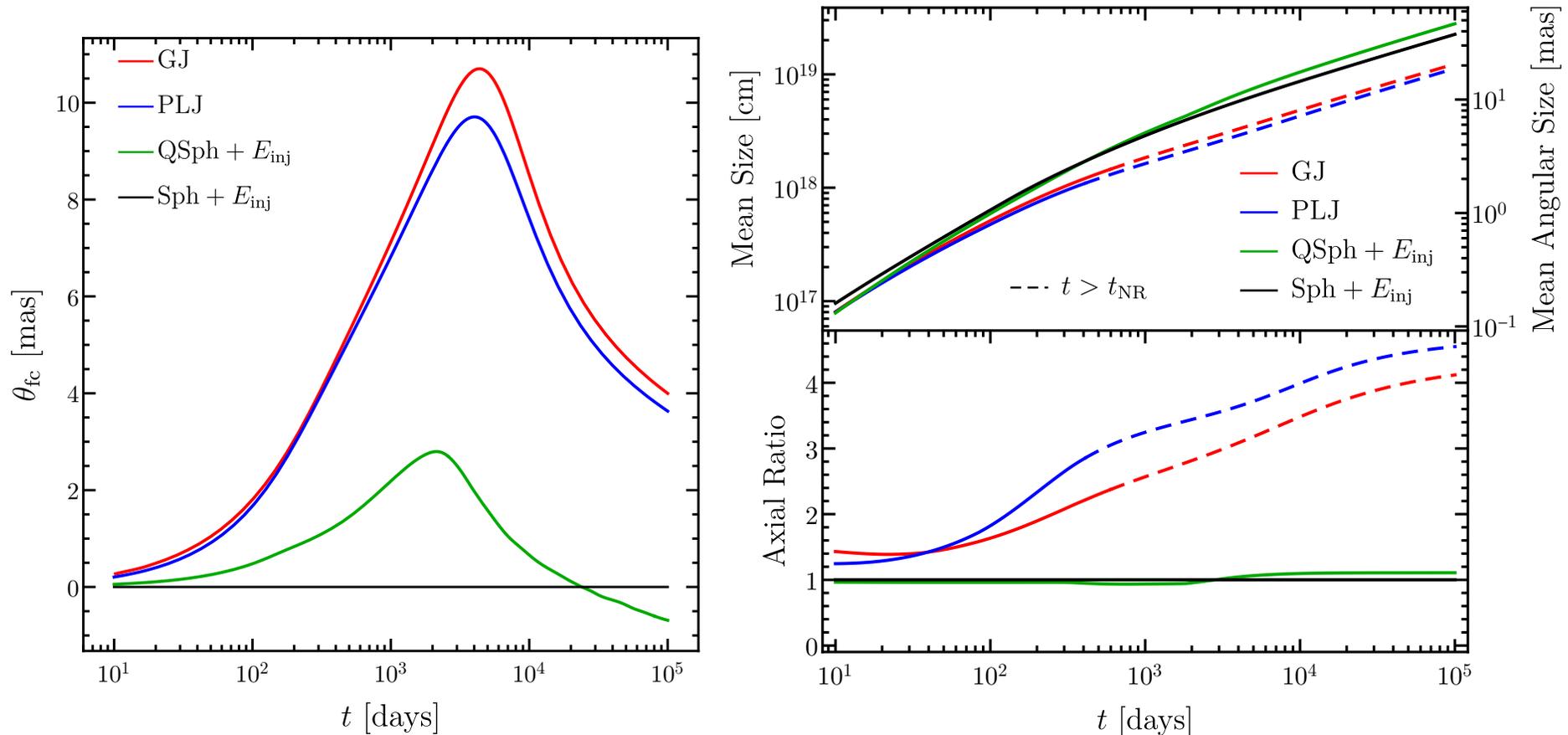
Outflow Structure: Breaking the Degeneracy (JG & Gill 18)

- The jet models decay faster (closer to post-peak data: $F_\nu \propto t^{-2.2}$)



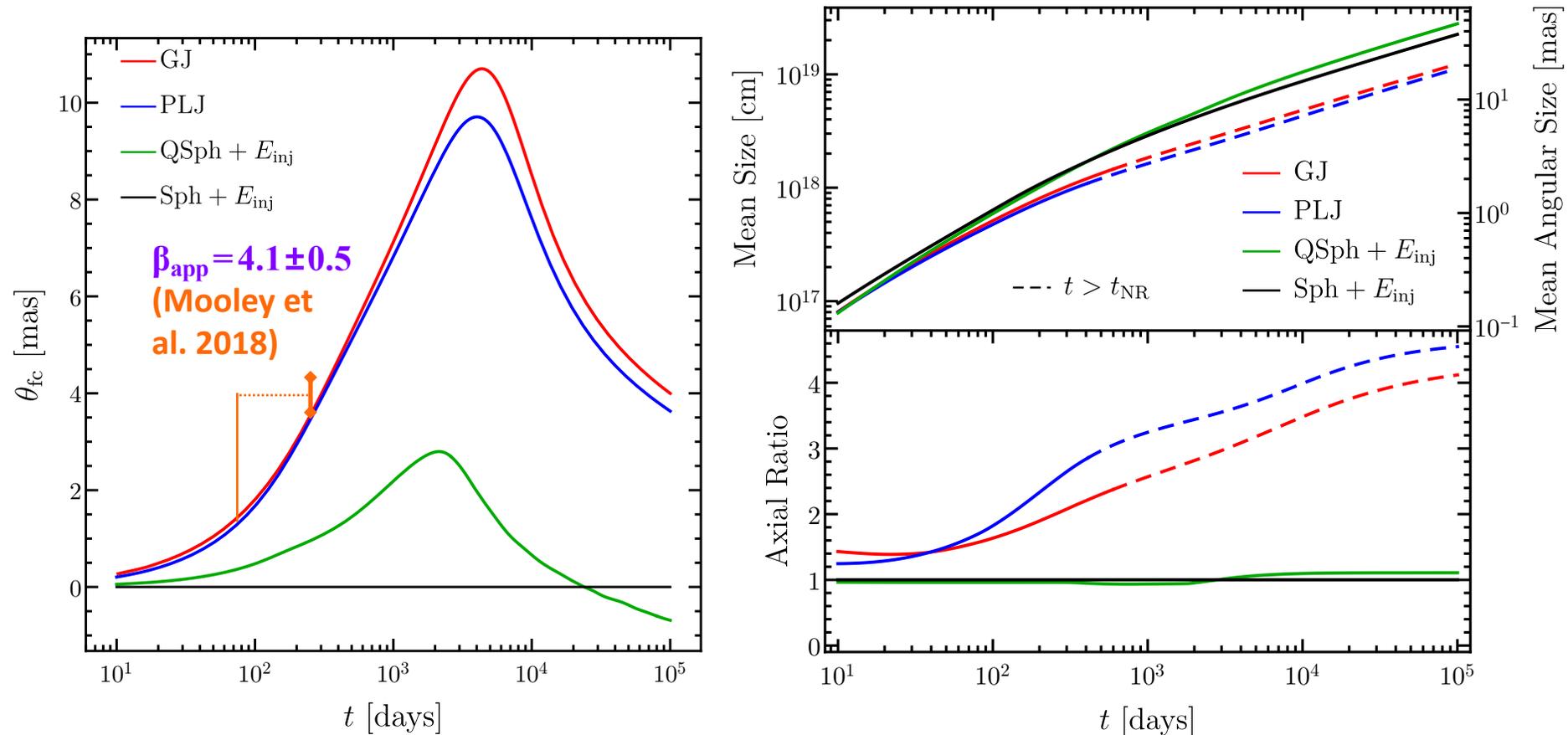
Afterglow Images: Flux Centroid, Size, Shape

- The flux centroid motion: a potentially powerful diagnostic
- It may be hard to tell apart models based on the image size alone, but a much higher axis-ratio is expected for jet models



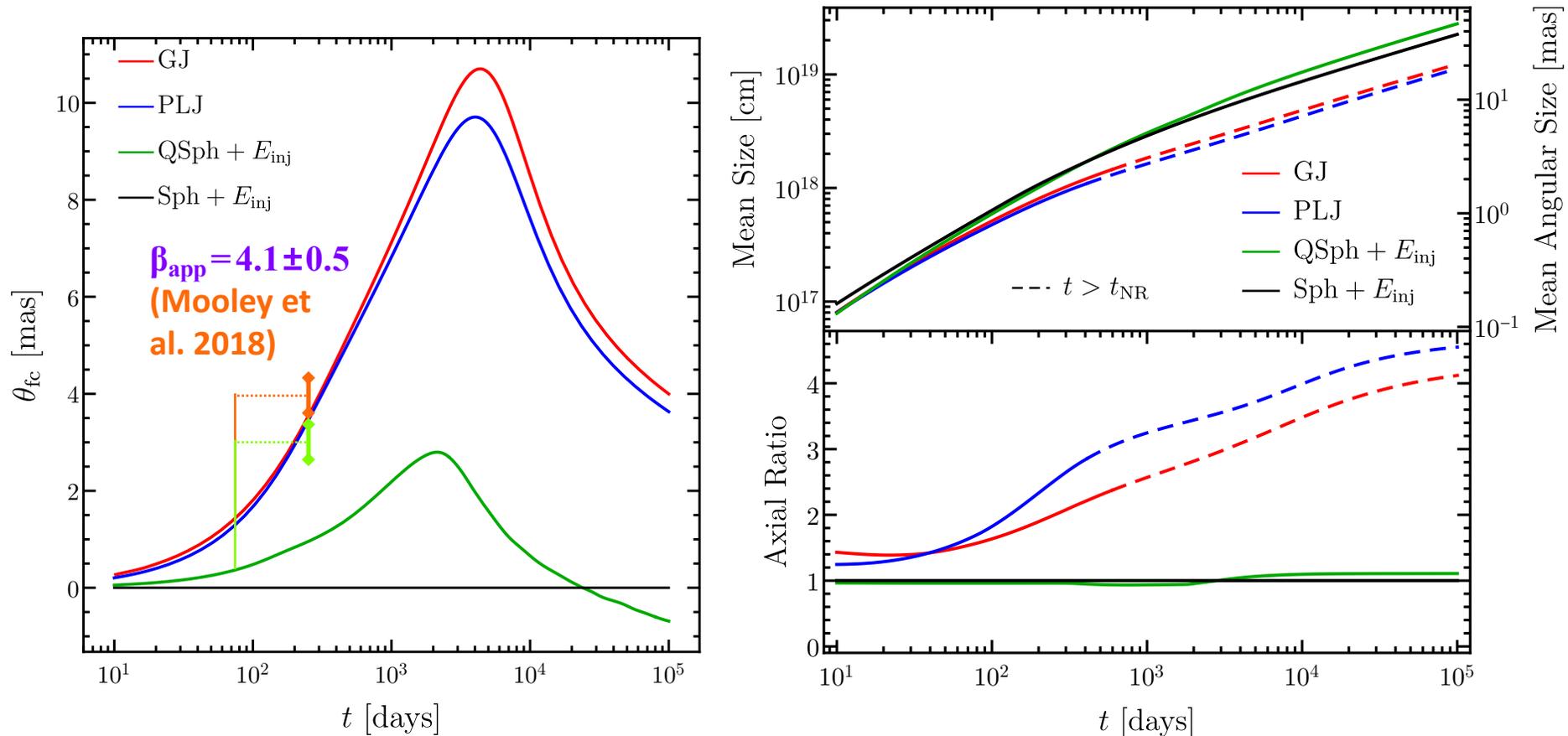
Afterglow Images: Flux Centroid, Size, Shape

- The flux centroid motion: a potentially powerful diagnostic
- It may be hard to tell apart models based on the image size alone, but a much higher axis-ratio is expected for jet models



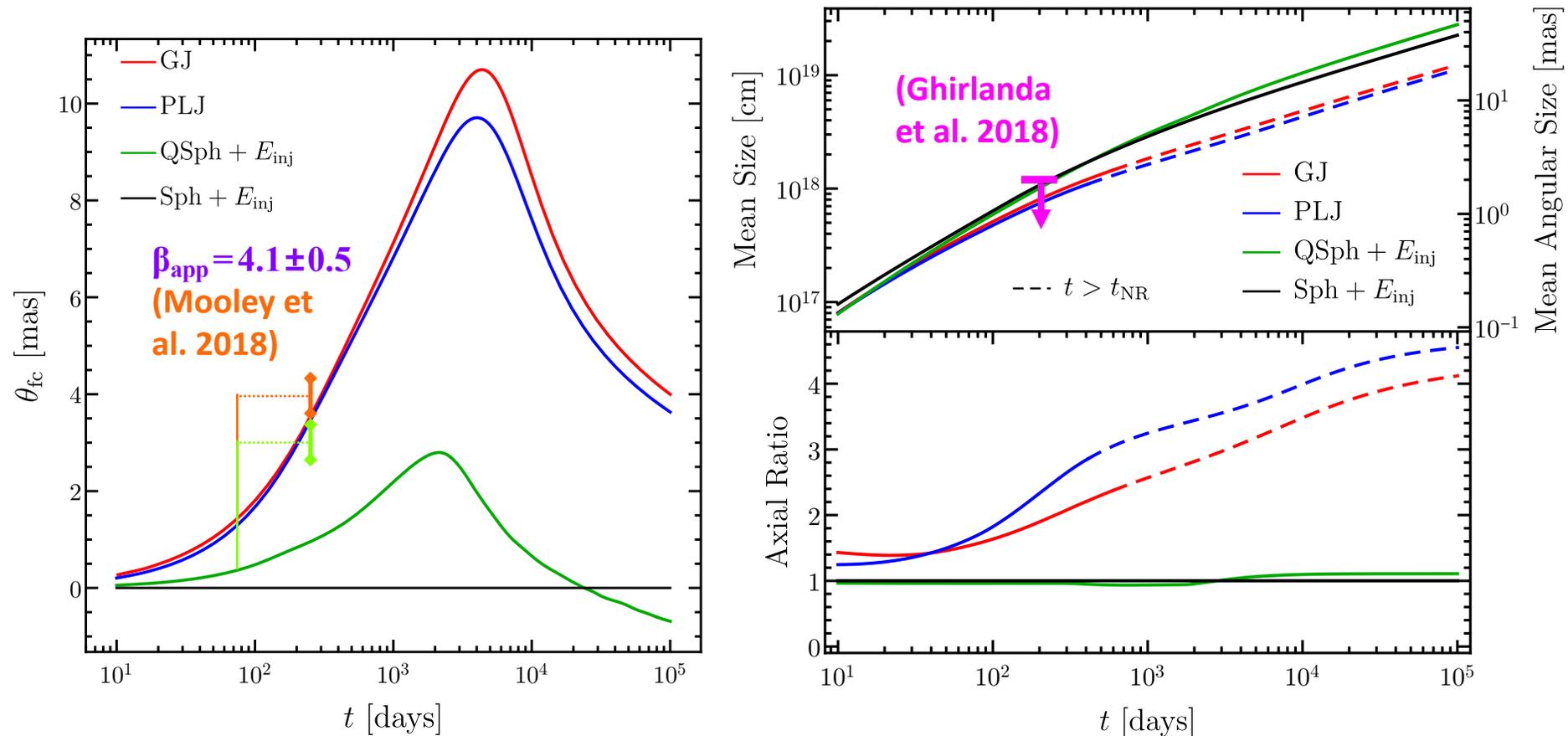
Afterglow Images: Flux Centroid, Size, Shape

- The flux centroid motion: a potentially powerful diagnostic
- It may be hard to tell apart models based on the image size alone, but a much higher axis-ratio is expected for jet models



Afterglow Images: Flux Centroid, Size, Shape

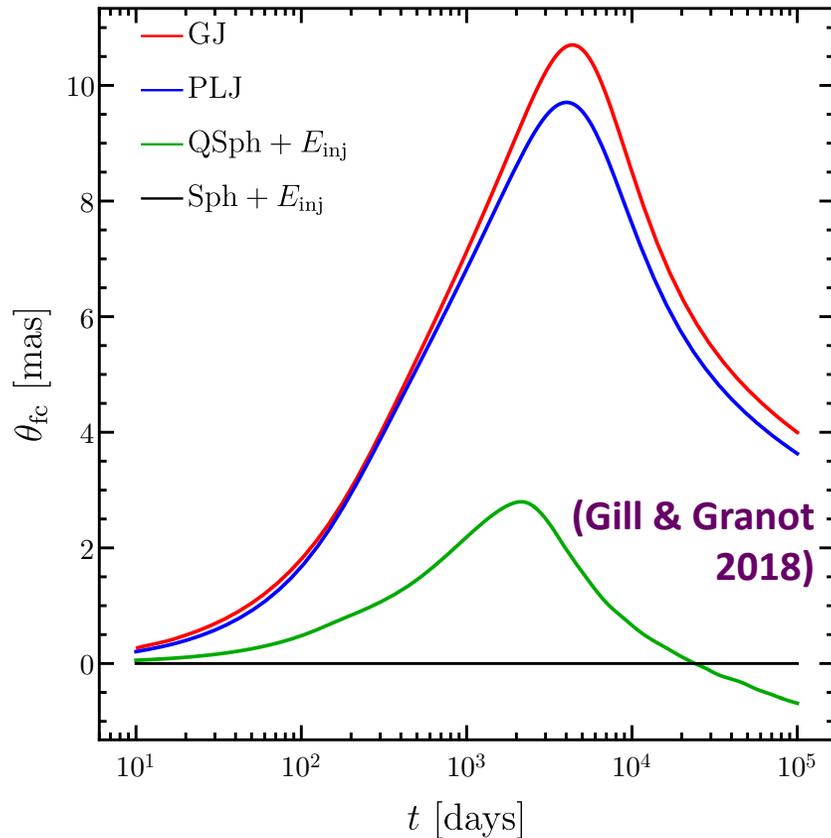
- The flux centroid motion: a potentially powerful diagnostic
- It may be hard to tell apart models based on the image size alone, but a much higher axis-ratio is expected for jet models



Afterglow Images: Flux Centroid, Size, Shape

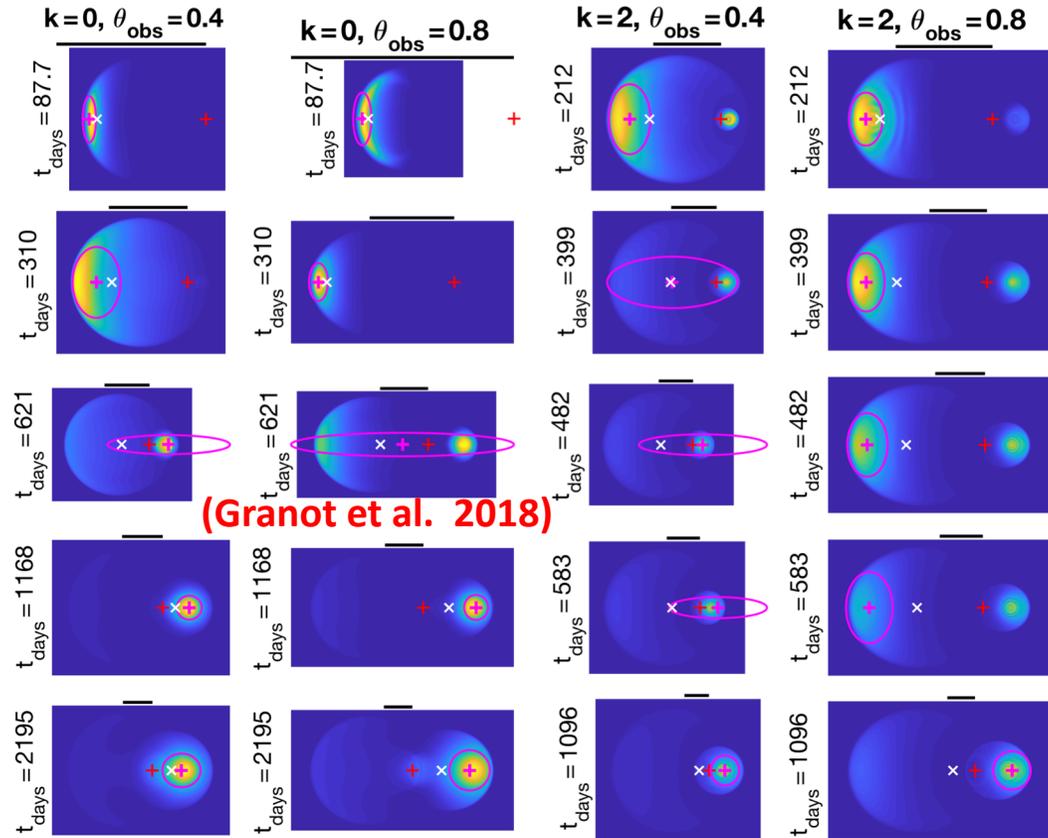
- The flux centroid motion: a potentially powerful diagnostic
- It may be hard to tell apart models based on the image size alone, but a much higher axis-ratio is expected for jet models

Radio flux centroid motion: semi-analytic



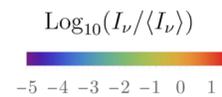
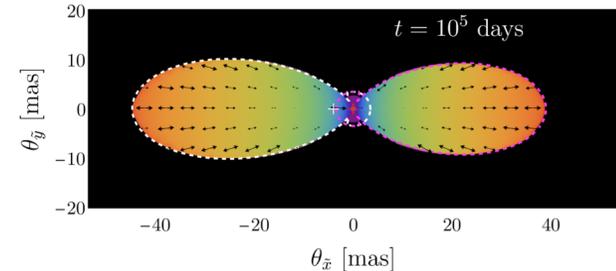
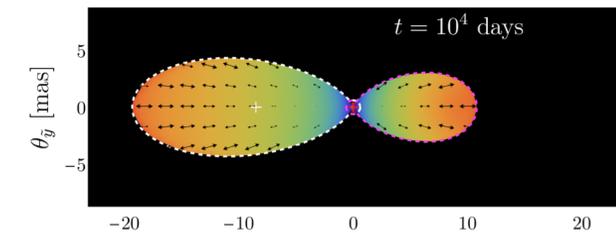
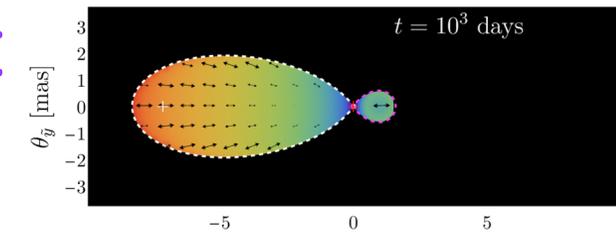
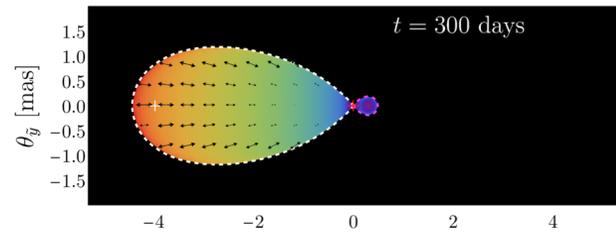
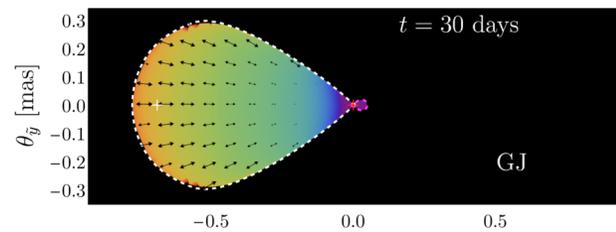
(Gill & JG 2018)

Agree with radio afterglow images from simulations



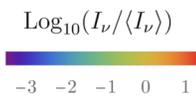
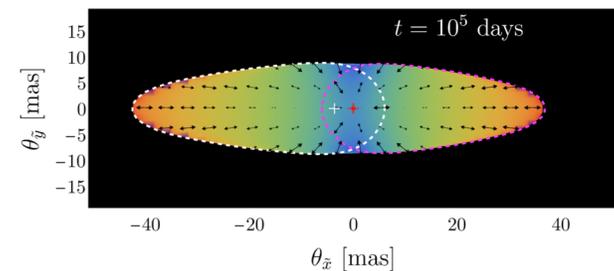
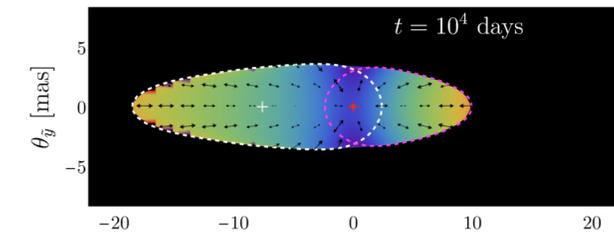
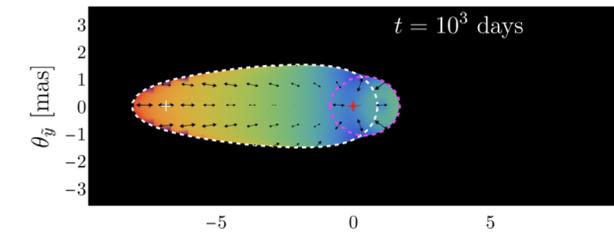
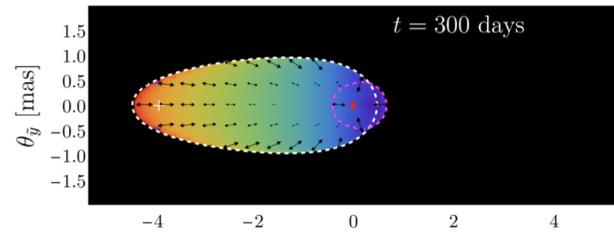
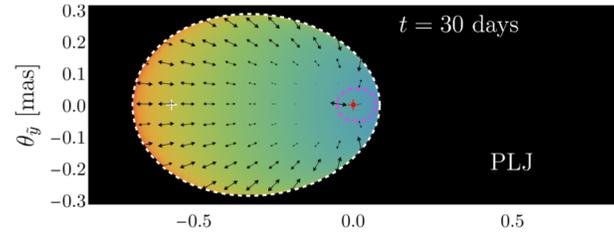
(JG, De Colle & Ramirez-Ruiz 2018)

Gaussian
Jet



Afterglow Images:

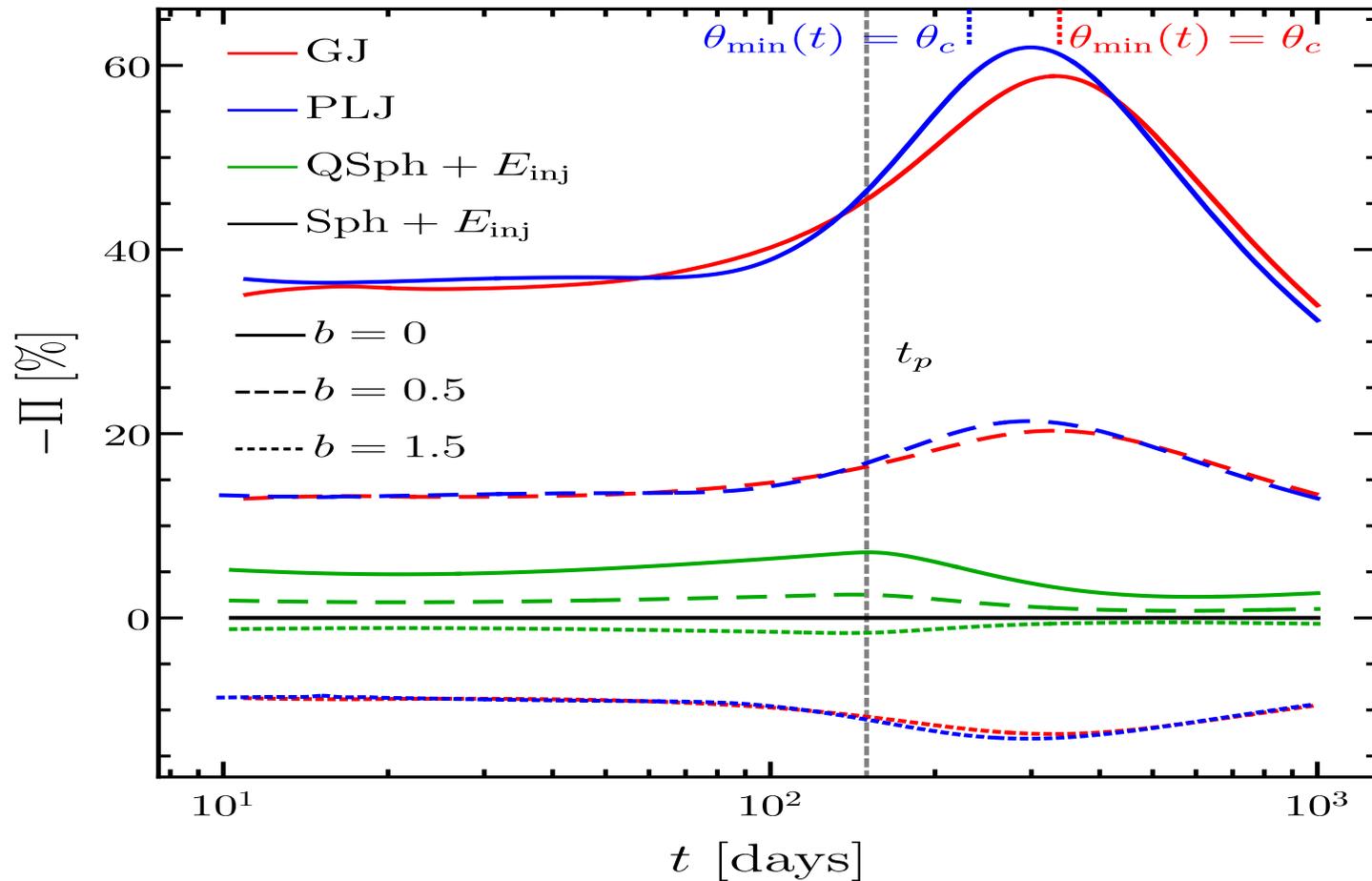
with a
polarization map



Power-
Law Jet

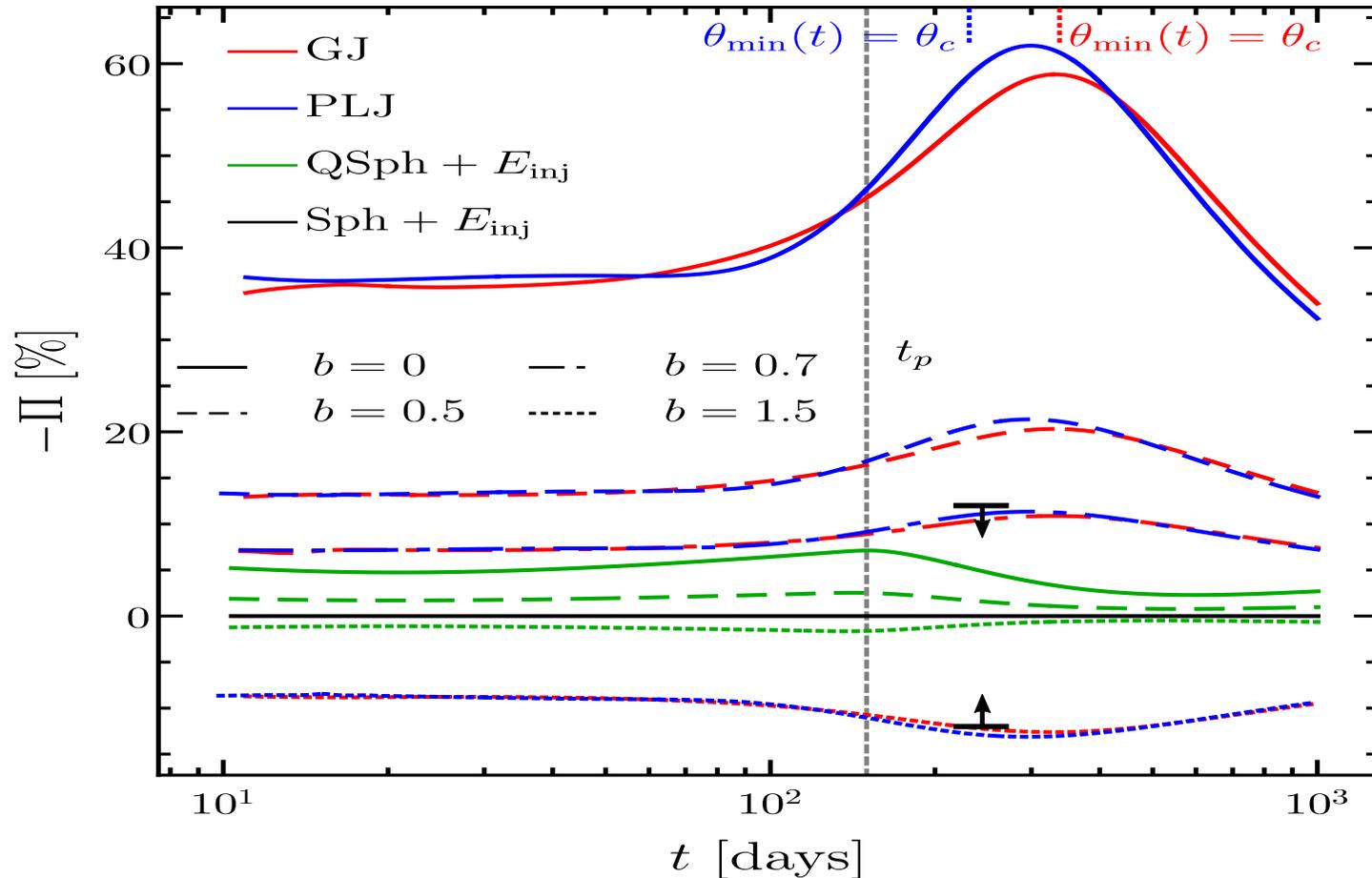
GRB 170817A: polarization UL \Rightarrow post-shock B-field

- Jet angular structure & θ_{obs} well constrained \Rightarrow breaks degeneracies
- Assuming a shock-produced B-field with $b \equiv 2\langle B_{\parallel}^2 \rangle / \langle B_{\perp}^2 \rangle$ (JG & König 03; Gill & JG 18)



GRB 170817A: polarization UL \Rightarrow post-shock B-field

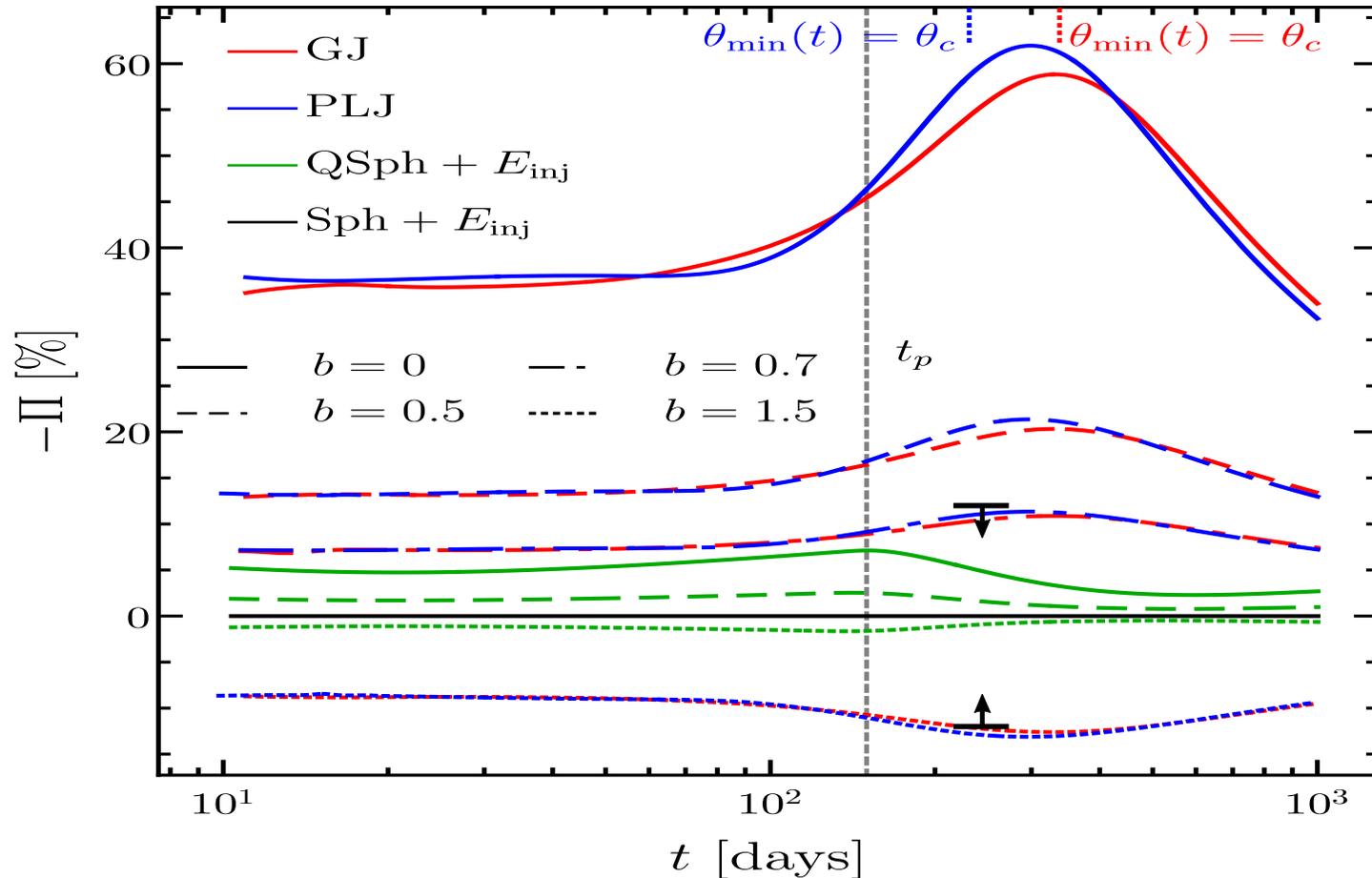
- Jet angular structure & θ_{obs} well constrained \Rightarrow breaks degeneracies
- Assuming a shock-produced B-field with $b \equiv 2\langle B_{\parallel}^2 \rangle / \langle B_{\perp}^2 \rangle$ (JG & Königl 03; Gill & JG 18)



Later: upper limit
 $P_{\text{lin}} < 12\%$ @
 $\nu = 2.8 \text{ GHz}$,
 $t = 244 \text{ days}$
 (Corsi + 2018)

GRB 170817A: polarization UL \Rightarrow post-shock B-field

- Jet angular structure & θ_{obs} well constrained \Rightarrow breaks degeneracies
- Assuming a shock-produced B-field with $b \equiv 2\langle B_{\parallel}^2 \rangle / \langle B_{\perp}^2 \rangle$ (JG & König 03; Gill & JG 18)



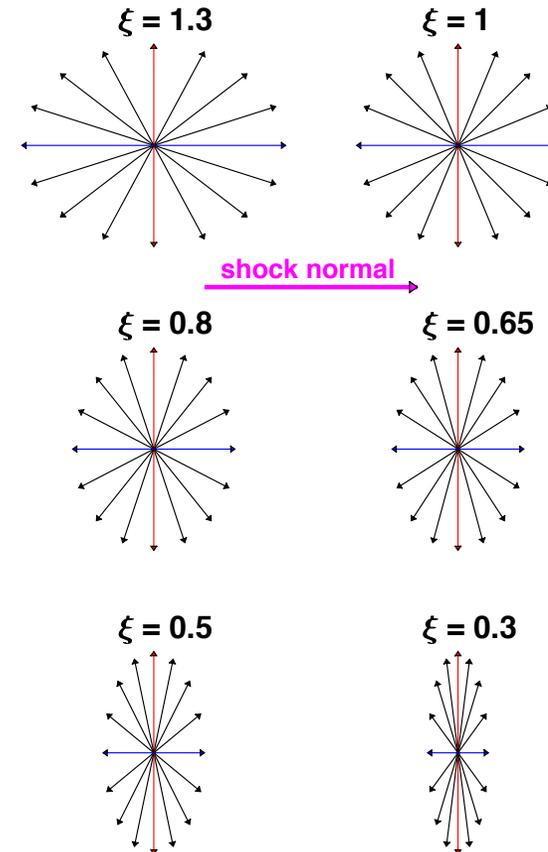
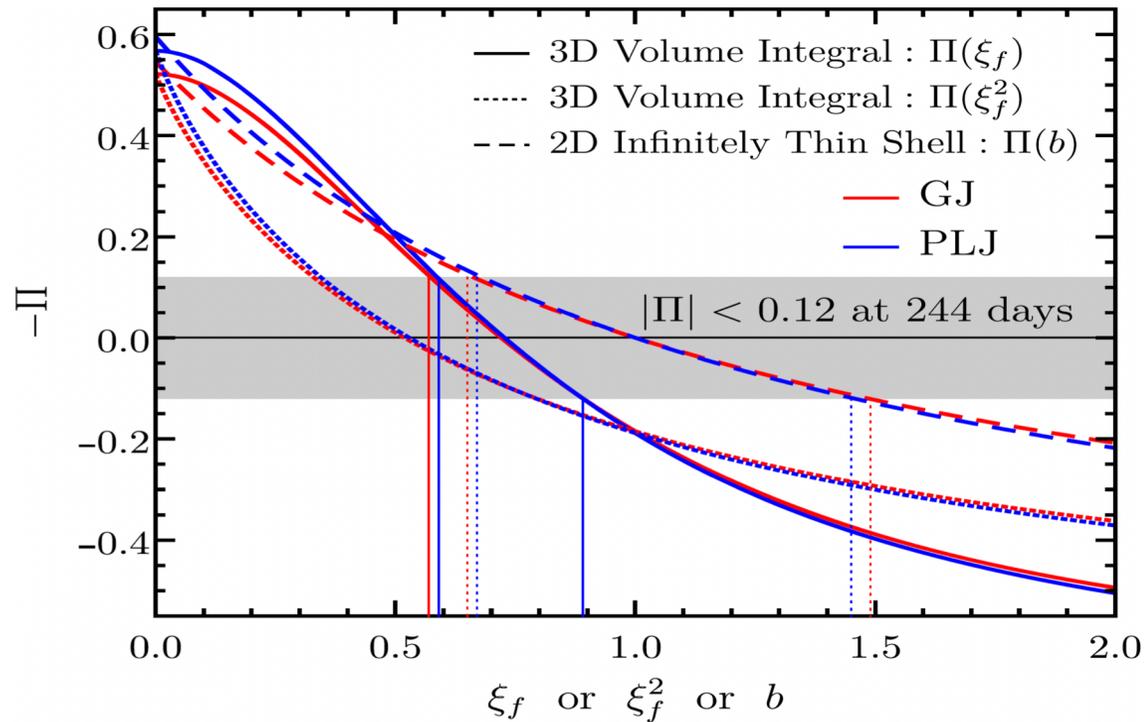
$0.7 \lesssim b \lesssim 1.5$
for jet models

Later: upper limit
 $P_{\text{lin}} < 12\%$ @
 $\nu = 2.8 \text{ GHz}$,
 $t = 244 \text{ days}$
(Corsi + 2018)

GRB 170817A: polarization UL \Rightarrow post-shock B-field

More realistic assumptions \Rightarrow B-field in collisionless shocks: (Gill & JG 2020)

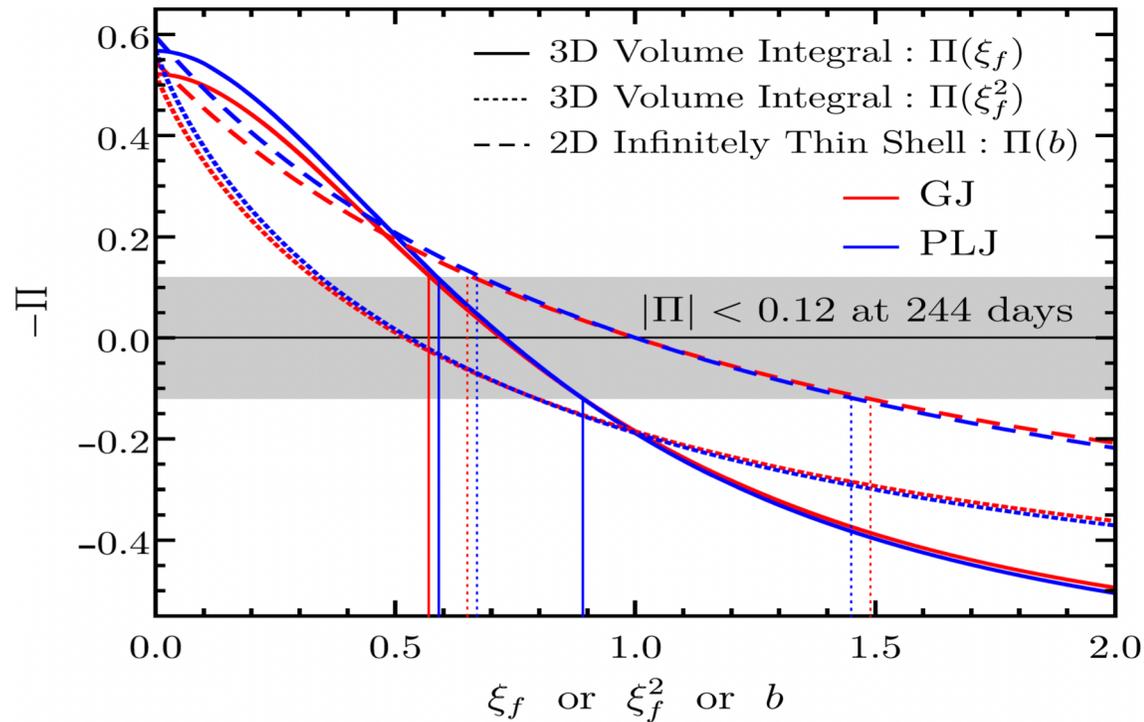
- 2D emitting shell \rightarrow 3D emitting volume (local BM76 radial profile)
- B-field evolution by faster radial expansion: $L'_r / L'_{\theta,\phi} \propto \chi^{(7-2k)/(8-2k)}$
- B-field isotropic in 3D with $B'_r \rightarrow \xi B'_r$ (Sari 1999); $\xi = \xi_f \chi^{(7-2k)/(8-2k)}$



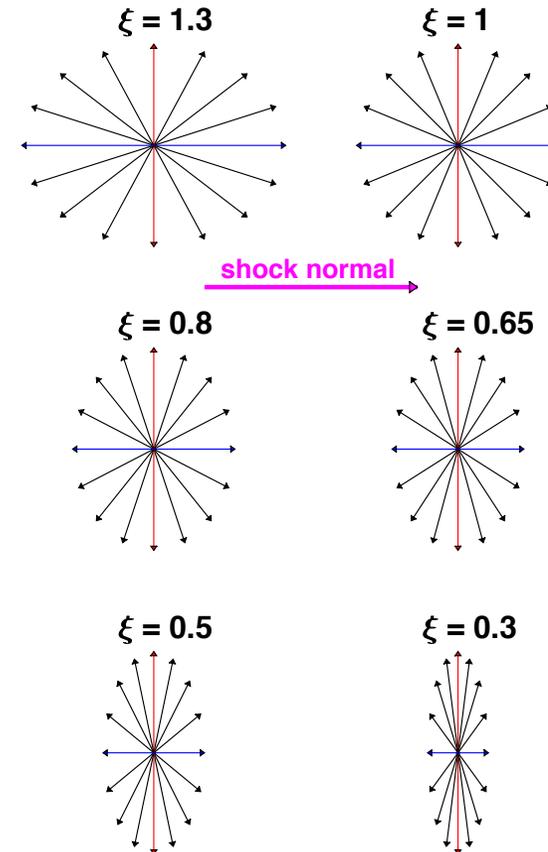
GRB 170817A: polarization UL \Rightarrow post-shock B-field

More realistic assumptions \Rightarrow B-field in collisionless shocks: (Gill & JG 2020)

- 2D emitting shell \rightarrow 3D emitting volume (local BM76 radial profile)
- B-field evolution by faster radial expansion: $L'_r / L'_{\theta,\phi} \propto \chi^{(7-2k)/(8-2k)}$
- B-field isotropic in 3D with $B'_r \rightarrow \xi B'_r$ (Sari 1999); $\xi = \xi_f \chi^{(7-2k)/(8-2k)}$



$$0.57 \lesssim \xi_f \lesssim 0.89$$

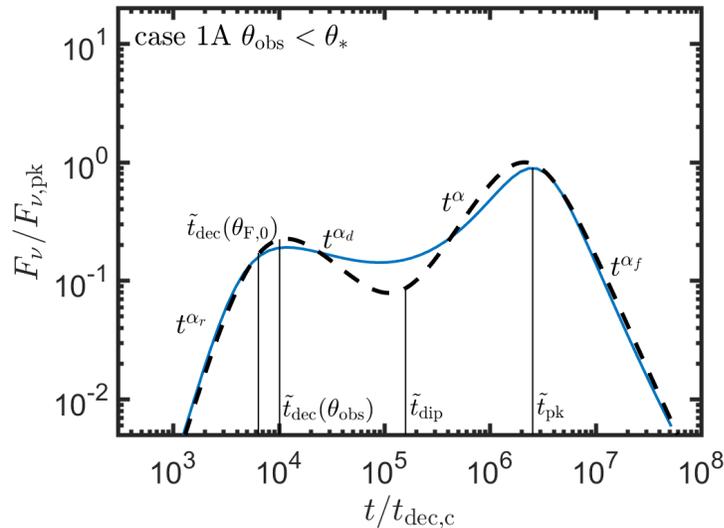


Predicted Off-Axis Lightcurves from Structured Jets

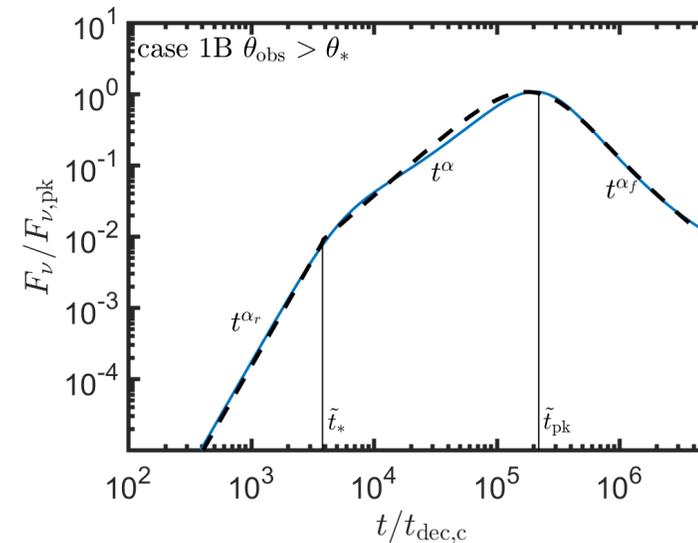
(Beniamini, JG & 2020; Beniamini, Gill & JG 2022)

- A general investigation of **Power-Law (+Gaussian) Jets**
- Provide detailed **analytic lightcurves**
- We find two main lightcurve types: **double** or **single** peaked

Double peaked LC: $\theta_{\text{obs}} < \theta_*$



Single peaked LC: $\theta_{\text{obs}} > \theta_*$



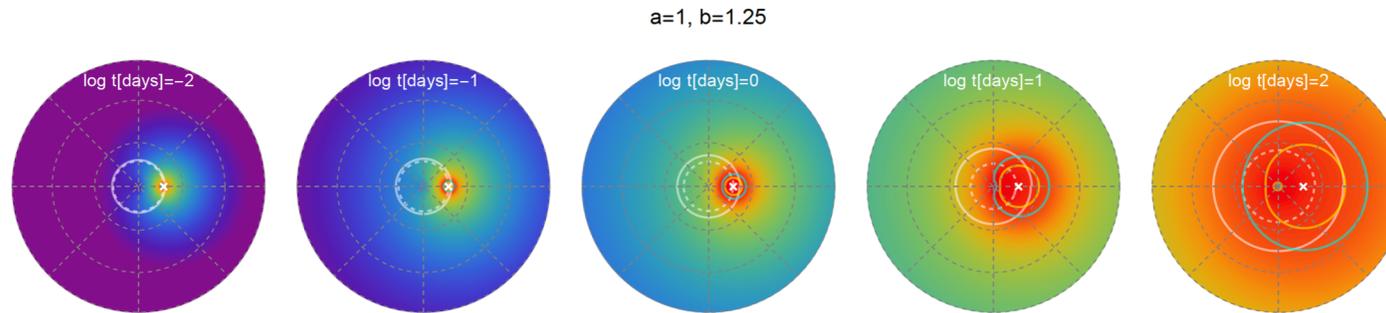
$$\theta_* \Gamma_0(\theta_*) = 1$$

Predicted Off-Axis Lightcurves from Structured Jets

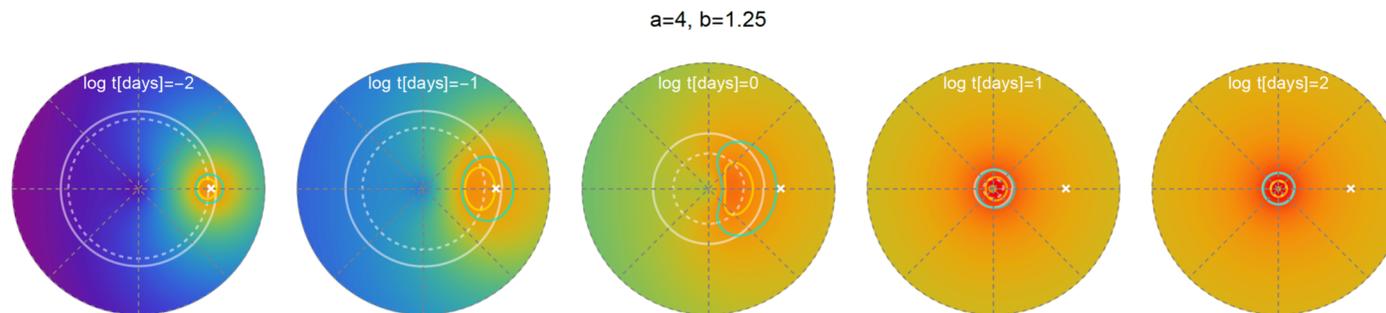
(Beniamini, JG & 2020; Beniamini, Gill & JG 2022)

- Map the most relevant parameter space from simulations of long / short GRB jets breaking out of the star / merger ejecta
 - ◆ ⇒ Consider different external density profiles
- Consider both shallow & steep jet angular profiles

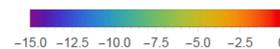
Shallow Jet:



Steep Jet:



$\log_{10}[dF_{\nu}/d\Omega]$

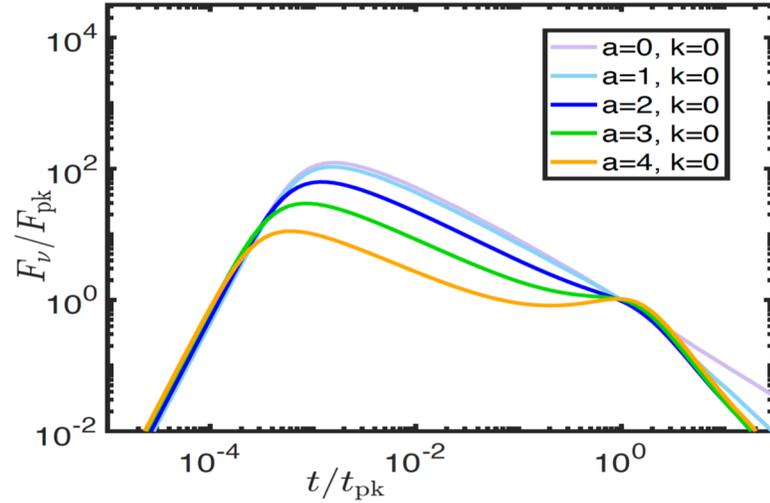


Screenshot

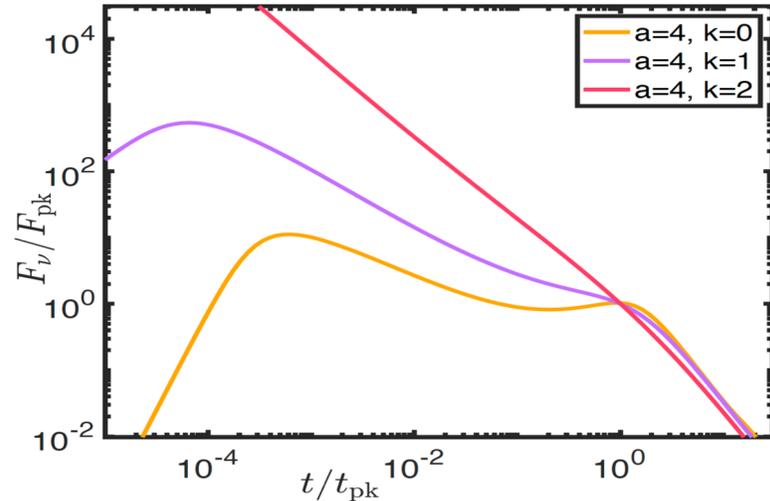
Predicted Off-Axis Lightcurves from Structured Jets

(Beniamini, JG & 2020; Beniamini, Gill & JG 2022)

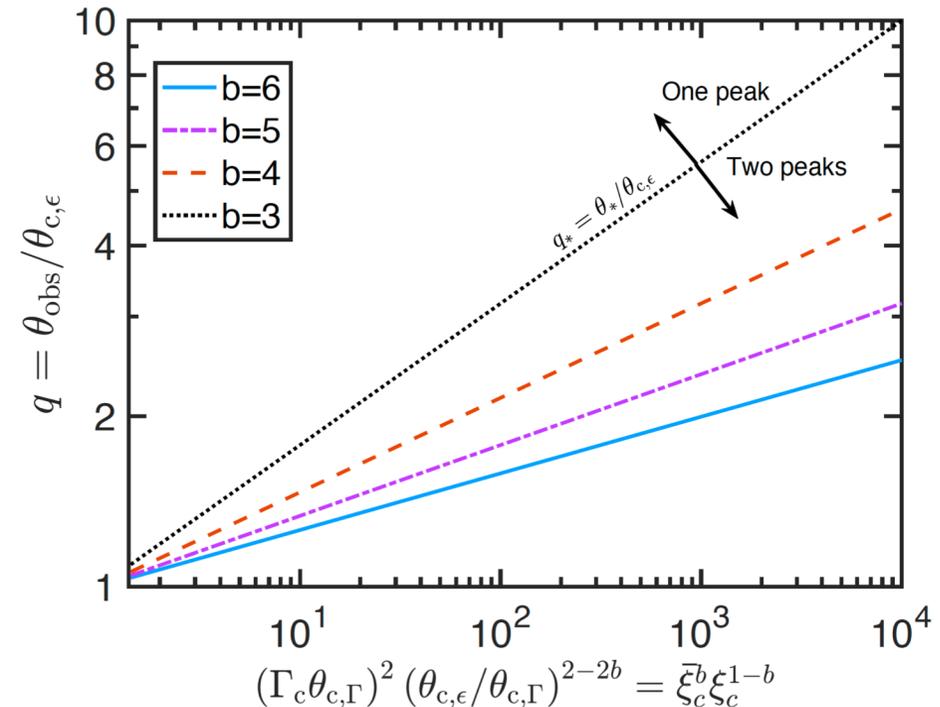
Shallow vs. Steep Jet:



External Density $\propto R^{-k}$:

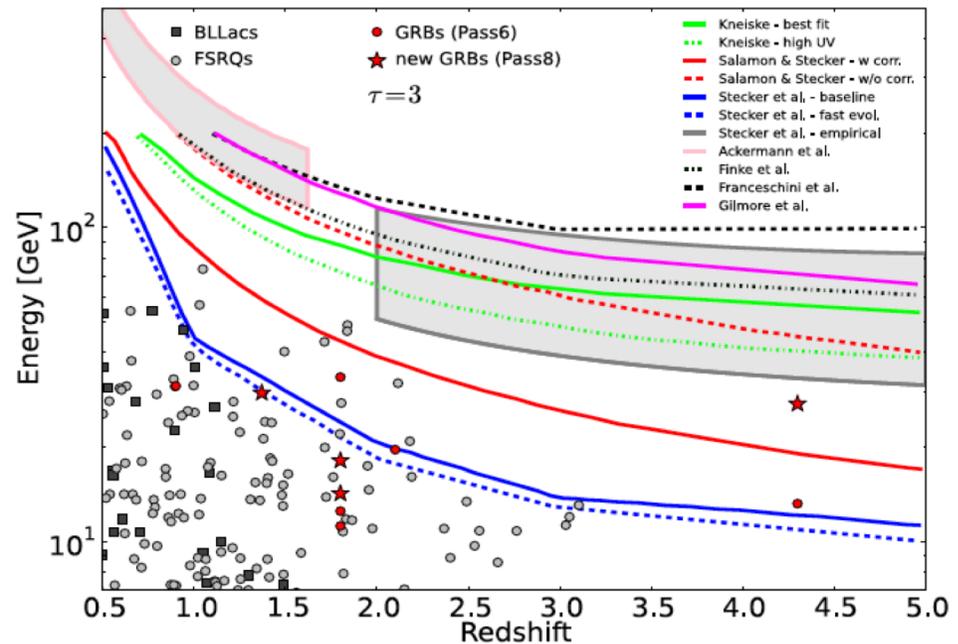
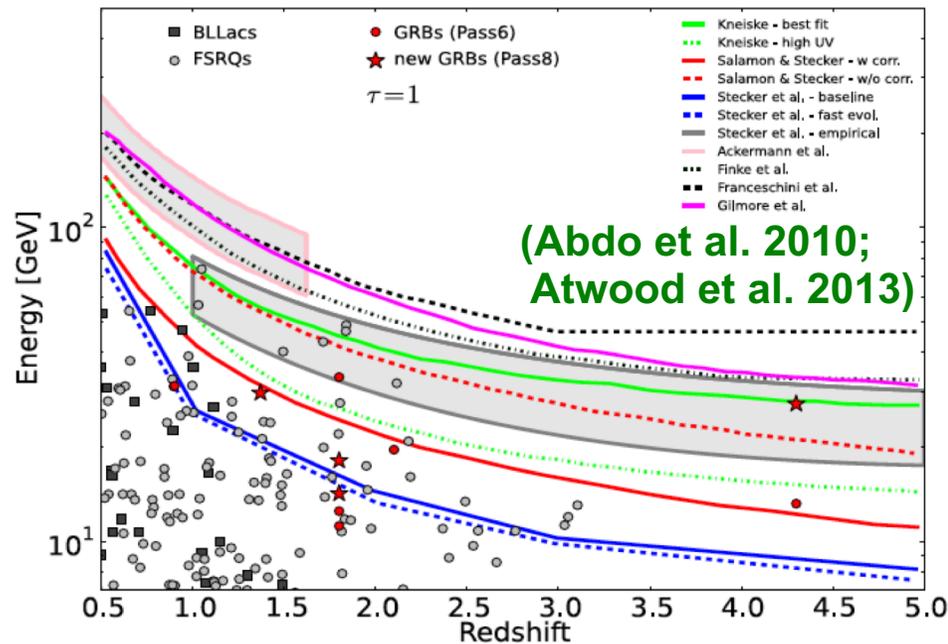


Single vs. Double Peaked LC:



Constraining the Opacity of the Universe

- γ -rays from distant sources can pair produce ($\gamma\gamma \rightarrow e^+e^-$) on the way to us with the extragalactic background light (EBL)
- This can test the transparency of the Universe and constrain EBL models (or the massive star formation rate at $z \gtrsim 1$)
- GRBs are already competitive with AGN, & probe higher z
- EBL likely detected (with blazars: LAT+IACTs; Dominguez+13; Acciari+19)

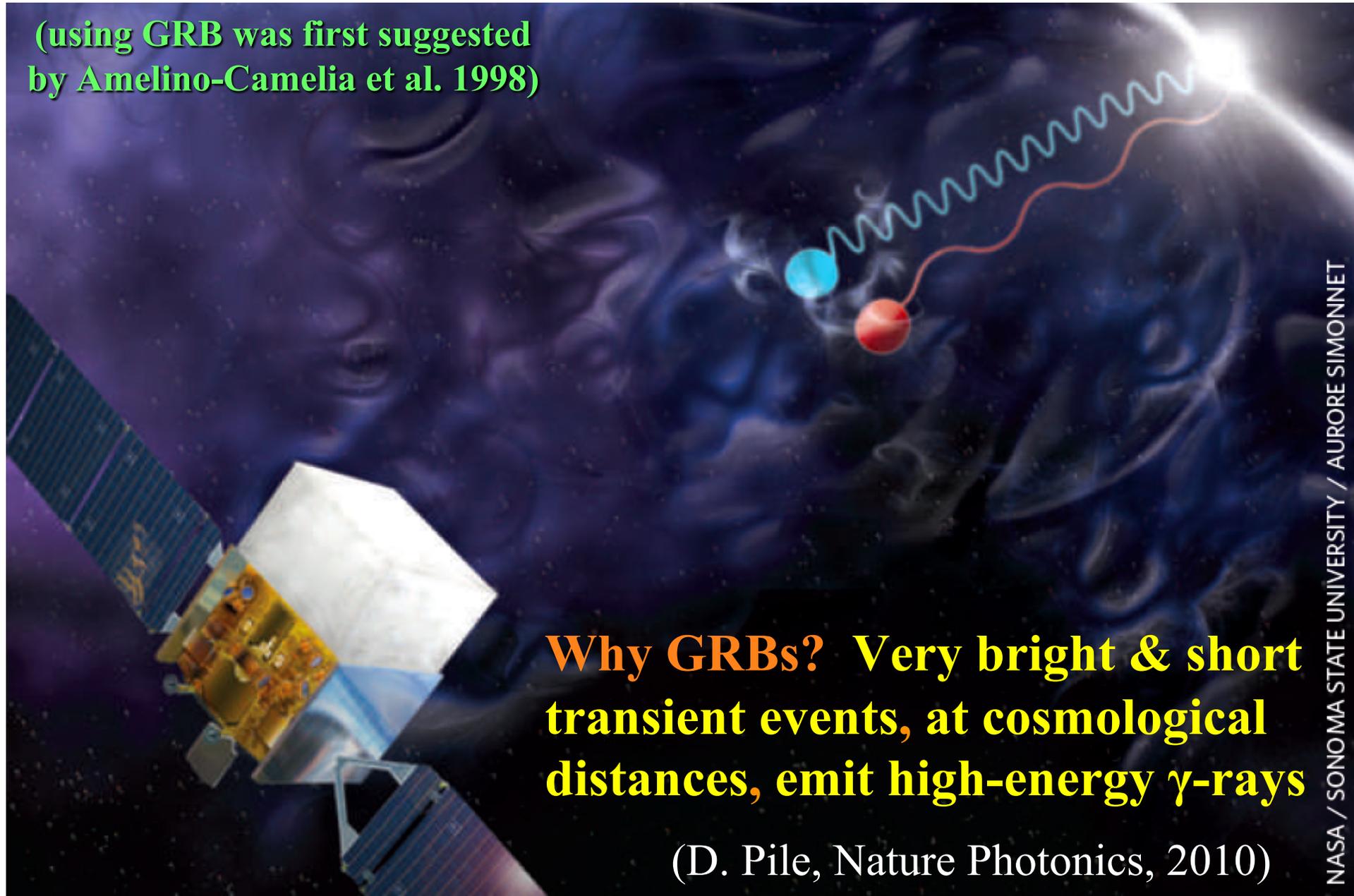


Testing for Lorentz Invariance Violation

(using GRB was first suggested
by Amelino-Camelia et al. 1998)

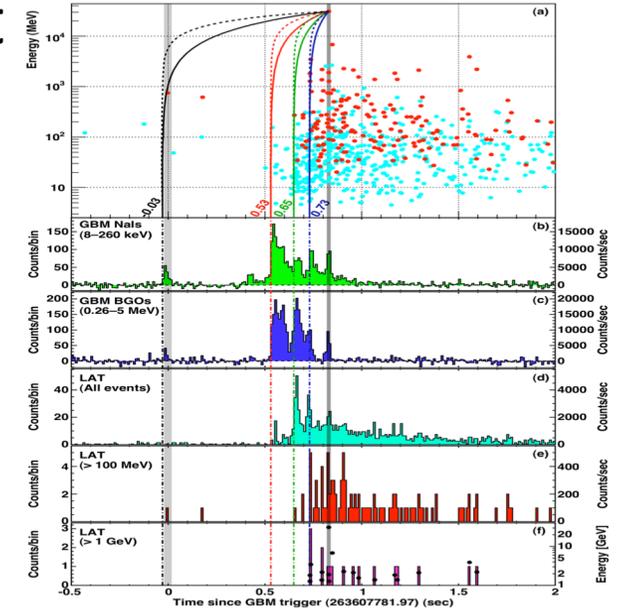
**Why GRBs? Very bright & short
transient events, at cosmological
distances, emit high-energy γ -rays**

(D. Pile, Nature Photonics, 2010)



Testing for Lorentz Invariance Violation

- GRB 090510 is much better than the rest (short, hard, very fine time structure)

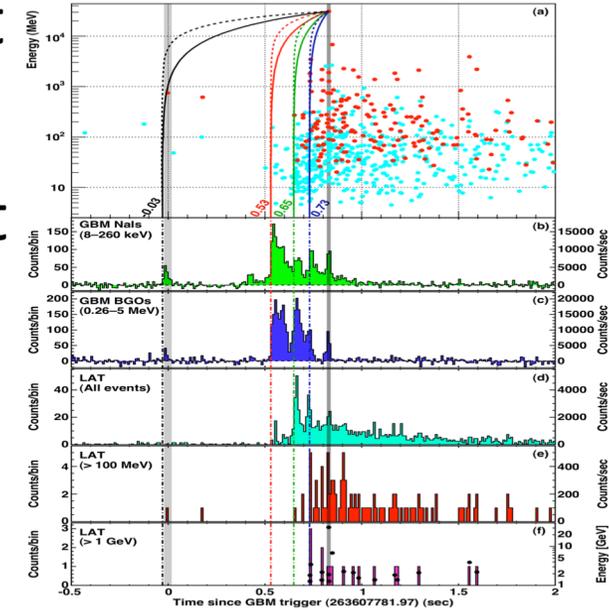


Testing for Lorentz Invariance Violation

- GRB 090510 is much better than the rest (short, hard, very fine time structure)
- Abdo+ 2009, *Nature*, 462, 331: 1st direct time-of-flight limit beyond Plank scale on linear ($n = 1$) energy dispersion:

$$v_{\text{ph}} / c \approx 1 \pm \frac{1}{2} (1+n) \left(E_{\text{ph}} / E_{\text{QG},n} \right)^n \quad E_{\text{QG},1} > 1.2 E_{\text{Planck}}$$

(robust, conservative, 2 independent methods)



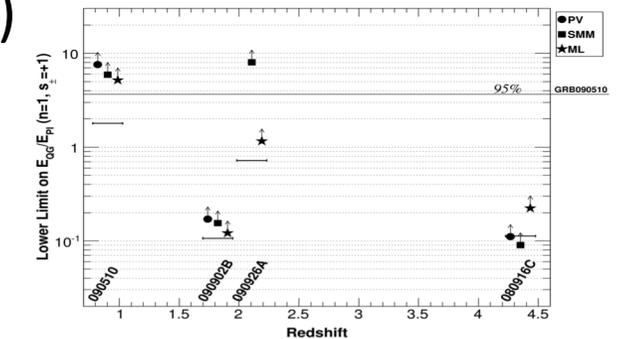
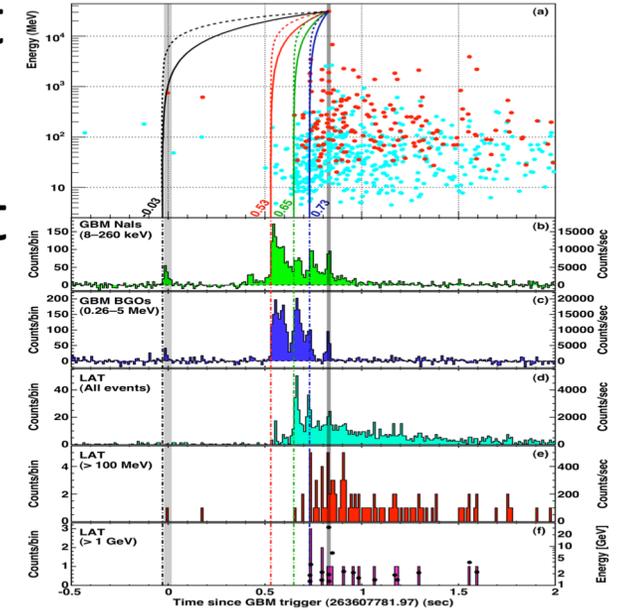
Testing for Lorentz Invariance Violation

- GRB 090510 is much better than the rest (short, hard, very fine time structure)
- Abdo+ 2009, *Nature*, 462, 331: 1st direct time-of-flight limit beyond Plank scale on linear ($n = 1$) energy dispersion:

$$v_{\text{ph}} / c \approx 1 \pm \frac{1}{2}(1+n) \left(E_{\text{ph}} / E_{\text{QG},n} \right)^n \quad E_{\text{QG},1} > 1.2 E_{\text{Planck}}$$

(robust, conservative, 2 independent methods)

- Vasileiou+ 2013: 3 different methods, 4 GRBs (090510 is still the best by far), the limits improved by factors of a few



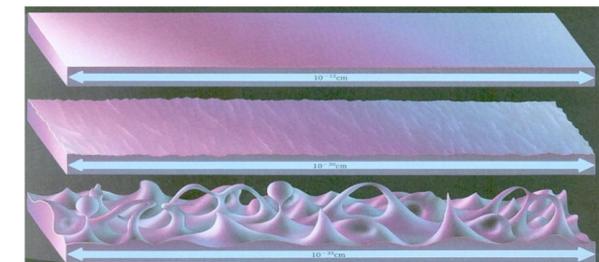
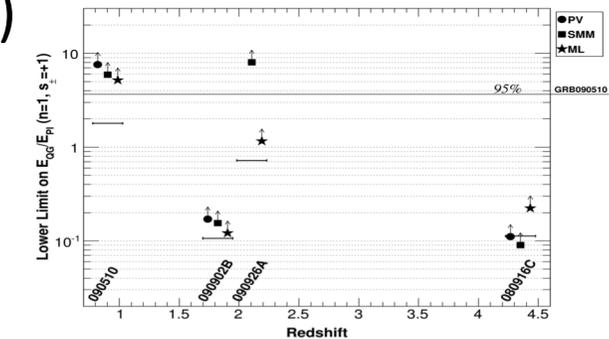
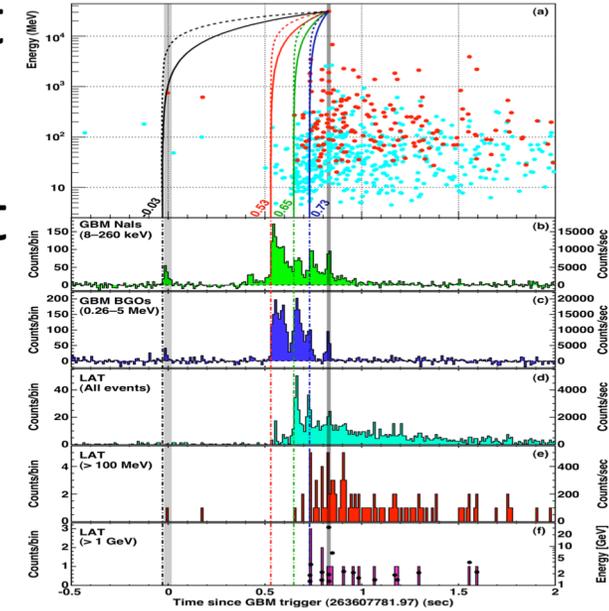
Testing for Lorentz Invariance Violation

- GRB 090510 is much better than the rest (short, hard, very fine time structure)
- Abdo+ 2009, *Nature*, 462, 331: 1st direct time-of-flight limit beyond Planck scale on linear ($n = 1$) energy dispersion:

$$v_{\text{ph}} / c \approx 1 \pm \frac{1}{2}(1+n) \left(E_{\text{ph}} / E_{\text{QG},n} \right)^n \quad E_{\text{QG},1} > 1.2 E_{\text{Planck}}$$

(robust, conservative, 2 independent methods)

- Vasileiou+ 2013: 3 different methods, 4 GRBs (090510 is still the best by far), the limits improved by factors of a few
- Vasileiou+ 2015, *Nature Phys.*, 11, 344: stochastic LIV – motivation: space-time foam (1st Planck-scale limit of its kind)



Conclusions: Short GRBs & Multimessenger Astrophysics

- GW170817 is unique with a wide range of implications
- GW speed: $\left| \frac{v_{GW}}{c} - 1 \right| \lesssim 4 \cdot 10^{-16}$; Kilonova: r-process elements
- Merger **Remnant**: BH or HMNS \rightarrow BH $\Rightarrow M_{TOV} \lesssim 2.17M_{\odot}$

Conclusions: Short GRBs & Multimessenger Astrophysics

- GW170817 is unique with a wide range of implications
- GW speed: $\left| \frac{v_{GW}}{c} - 1 \right| \lesssim 4 \cdot 10^{-16}$; Kilonova: r-process elements
- Merger **Remnant**: BH or HMNS \rightarrow BH $\Rightarrow M_{TOV} \lesssim 2.17M_{\odot}$
- Two main types of explanations for the rising afterglow flux energy distribution with proper velocity (r) or with angle (θ)

Conclusions: Short GRBs & Multimessenger Astrophysics

- GW170817 is unique with a wide range of implications
- GW speed: $\left| \frac{v_{GW}}{c} - 1 \right| \lesssim 4 \cdot 10^{-16}$; Kilonova: r-process elements
- Merger **Remnant**: BH or HMNS \rightarrow BH $\Rightarrow M_{TOV} \lesssim 2.17M_{\odot}$
- Two main types of explanations for the rising afterglow flux energy distribution with proper velocity (r) or with angle (θ)
- Possible diagnostics to distinguish between them
 - ◆ The post-peak flux decay slope
 - ◆ Flux centroid motion or image axis ratio

Conclusions: Short GRBs & Multimessenger Astrophysics

- GW170817 is unique with a wide range of implications
- GW speed: $\left| \frac{v_{GW}}{c} - 1 \right| \lesssim 4 \cdot 10^{-16}$; Kilonova: r-process elements
- Merger **Remnant**: BH or HMNS \rightarrow BH $\Rightarrow M_{TOV} \lesssim 2.17M_{\odot}$
- Two main types of explanations for the rising afterglow flux energy distribution with proper velocity (r) or with angle (θ)
- Possible diagnostics to distinguish between them
 - ◆ The post-peak flux decay slope
 - ◆ Flux centroid motion or image axis ratio
- Later flux centroid motion observations: $\beta_{app} = 4.1 \pm 0.5$

Conclusions: Short GRBs & Multimessenger Astrophysics

- GW170817 is unique with a wide range of implications
- GW speed: $\left| \frac{v_{GW}}{c} - 1 \right| \lesssim 4 \cdot 10^{-16}$; Kilonova: r-process elements
- Merger **Remnant**: BH or HMNS \rightarrow BH $\Rightarrow M_{TOV} \lesssim 2.17M_{\odot}$
- Two main types of explanations for the rising afterglow flux energy distribution with proper velocity (r) or with angle (θ)
- Possible diagnostics to distinguish between them
 - ◆ The post-peak flux decay slope
 - ◆ Flux centroid motion or image axis ratio
- Later flux centroid motion observations: $\beta_{app} = 4.1 \pm 0.5$
- Polarization UL: shock-produced B-field $0.57 \lesssim \xi_0 \lesssim 0.89$

Conclusions: Short GRBs & Multimessenger Astrophysics

- GW170817 is unique with a wide range of implications
- GW speed: $\left| \frac{v_{GW}}{c} - 1 \right| \lesssim 4 \cdot 10^{-16}$; Kilonova: r-process elements
- Merger **Remnant**: BH or HMNS \rightarrow BH $\Rightarrow M_{TOV} \lesssim 2.17M_{\odot}$
- Two main types of explanations for the rising afterglow flux energy distribution with proper velocity (r) or with angle (θ)
- Possible diagnostics to distinguish between them
 - ◆ The post-peak flux decay slope
 - ◆ Flux centroid motion or image axis ratio
- Later flux centroid motion observations: $\beta_{app} = 4.1 \pm 0.5$
- Polarization UL: shock-produced B-field $0.57 \lesssim \xi_0 \lesssim 0.89$
- Predicted **off-axis lightcurves** from **structured jets**

Conclusions: Short GRBs & Multimessenger Astrophysics

- GW170817 is unique with a wide range of implications
- GW speed: $\left| \frac{v_{GW}}{c} - 1 \right| \lesssim 4 \cdot 10^{-16}$; Kilonova: r-process elements
- Merger **Remnant**: BH or HMNS \rightarrow BH $\Rightarrow M_{TOV} \lesssim 2.17M_{\odot}$
- Two main types of explanations for the rising afterglow flux energy distribution with proper velocity (r) or with angle (θ)
- Possible diagnostics to distinguish between them
 - ◆ The post-peak flux decay slope
 - ◆ Flux centroid motion or image axis ratio
- Later flux centroid motion observations: $\beta_{app} = 4.1 \pm 0.5$
- Polarization UL: shock-produced B-field $0.57 \lesssim \xi_0 \lesssim 0.89$
- Predicted **off-axis lightcurves** from **structured jets**
- GRBs can also constrain Lorentz Invariance Violation or the EBL

The End